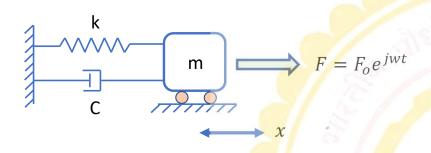
Lec 8 CONCEPT OF MECHANICAL IMPEDANCE



Harmonic Response: $\dot{x} = j\omega x = \omega x e^{\frac{\pi}{2}j}$

$$\ddot{x} = -\omega^2 x = \omega^2 x e^{\pi j}$$

Response: Displacement $x = x_0 e^{j(wt \cdot \varphi)}$ $= (x_0 e^{-j\varphi}) \cdot e^{jwt}$ Phase lag of displacement w.r.t force $= (x_0 e^{-j\varphi}) \cdot e^{jwt}$ Include phase information

Velocity leads displacement by $\frac{\pi}{2}$ & acceleration leads velocity by $\frac{\pi}{2}$

Equation of motion: $m\ddot{x} + c\dot{x} + kx = F = F_o e^{j\omega t}$

$$(-m\omega^2 + j\omega C + k)x = F$$

Effective dynamic stiffness: $\frac{F}{x} = -m\omega^2 + j\omega C + k$

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Special case: Negligible inertial force

$$(k_d)_{eff} = k + j\omega C$$

$$= (k + k^*j)$$

$$= \frac{A}{L}(\gamma + \gamma^*j)$$

$$F = F_o e^{jwt}$$
 Area, A

Define $\eta = \frac{\gamma^*}{\gamma}$ Mechanical loss factor

$$(k_d)_{eff} = \frac{A\gamma}{L}(1 + \eta j)$$

Complex Young Modulus

$$\bar{\gamma} = \gamma (1 + \eta j)$$

Under dynamic conditions

$$|\gamma|_{dynamic} = \gamma \sqrt{1 + \eta^2} > |\gamma|_{static}$$

Exact relation between η , ξ , and C:

$$k + j\omega C \equiv k + k\eta j$$
$$\omega C = k\eta$$
$$\eta = \left(\frac{C}{k}\right)\omega$$

Note:
$$\omega = \omega_N$$

 $[k] = \left\lceil \frac{N}{m} \right\rceil$

$$C = 2m\omega \xi$$

 $C_c = Crtical damping$

$$\omega^2 = \frac{k}{m}$$

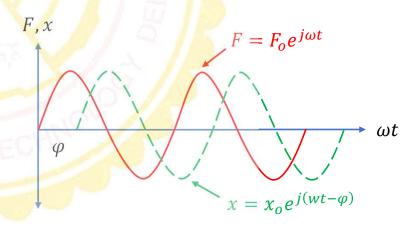
On Solving $\eta = 2\xi$

$$F = (-m\omega^2 + j\omega\mathcal{C} + k)x$$

If force, $F = F_0 e^{j\omega t}$

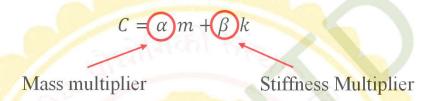
Displacement, $x = x_0 e^{j(wt - \varphi)}$

If no damping, $\varphi = 0$



 $[C] = \left[\frac{Ns}{m}\right]$

Rayleigh Damping



How $\alpha \& \beta$ related to $\xi \& \eta$?

Solution 1:

$$C = \alpha m + \beta k = \frac{k\eta}{\omega}$$

$$\Rightarrow \alpha = 0, \qquad \beta = \frac{\eta}{\omega} = \frac{2\xi}{\omega} = \frac{\xi}{\pi f}$$

Solution 2: Preferred in case of frequency range

Divide both sides by $2m\omega$

$$\alpha m + \beta k = C = 2mw\xi$$

$$\frac{\alpha m}{2m\omega} + \frac{\beta(m\omega^2)}{2m\omega} = \xi$$

$$\frac{\alpha}{2\omega} + \frac{\beta\omega}{2} = \xi$$

Infinite solutions

$$\frac{\alpha}{4\pi f} + \beta \pi f = \xi$$

Let frequency range be (f_1, f_2)

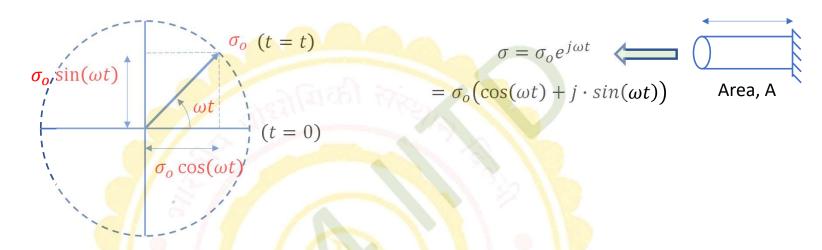
Two possible equations

$$\frac{\alpha}{4\pi f_1} + \beta \pi f_1 = \xi$$
 ... (1)

$$\frac{\alpha}{4\pi f_1} + \beta \pi f_1 = \xi \qquad \dots (1)$$

$$\frac{\alpha}{4\pi f_2} + \beta \pi f_2 = \xi \qquad \dots (2)$$

Solve these 2 equations to get $(\alpha, \beta) \to \text{Representation of range } (f_1, f_2)$



Either of the real or the imaginary component may be chosen

Strain,
$$\epsilon = \frac{\sigma}{\bar{\gamma}} = \frac{\sigma_o e^{j\omega t}}{\gamma(1+\eta j)} \times \frac{(1-\eta j)}{(1-\eta j)} \qquad \eta \ll 1$$

$$= \frac{\sigma_o (1-\eta j) e^{j\omega t}}{\gamma(1+\eta^2)}$$

$$\approx \frac{\sigma_o}{\gamma} \cdot e^{-j\eta} \cdot e^{j\omega t}$$

$$\epsilon = \frac{\sigma_o}{\gamma} \cdot e^{j(\omega t - \eta)}$$
Phase lag

If $\sigma = \sigma_0 \cos(\omega t)$ \rightarrow Projection in X axis

Then $\epsilon = \epsilon_0 \cos(\omega t - \eta)$

Alternatively, if $\sigma = \sigma_o \sin(\omega t)$, $\epsilon = \epsilon_o \sin(\omega t - \eta)$

If $\eta = 0$, $\sigma \& \epsilon$ in phase:

$$\frac{\sigma_o}{\gamma(1+\eta^2)} \approx \frac{\sigma_o}{\gamma}$$

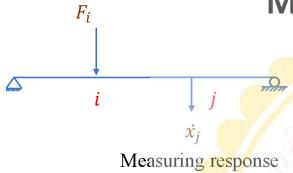
$$\in = \frac{\sigma_o}{\gamma} \cdot e^{\mathbf{j}(\omega t - \eta)}$$

$$= \left(\frac{\sigma_o}{\gamma} \cdot e^{-j\eta}\right) \cdot e^{j\omega t}$$

$$= (\epsilon_{o,r} + \epsilon_{o,i}) \cdot e^{j\omega t}$$

Phase information included in the magnitude part

Mechanical Impedance



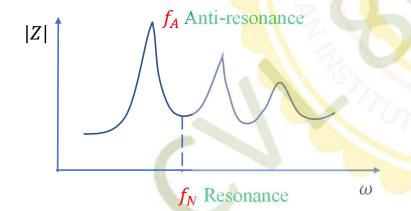
$$Z_{ij} = \frac{F_i}{\dot{x}_j}$$

Ratio of force to velocity

If
$$i = j$$
 then

$$Z_{ij} = \frac{F_i}{\dot{x}_i} \qquad \longrightarrow$$

Ratio of force applied at one point in the structure to the resulting velocity at same point in the direction of force.



Similar to FRF

$$F = F_o e^{jwt}$$

$$F = C\dot{x}$$

$$Z_{damper} = C$$

$$F = F_o e^{jwt}$$

$$F = kx = \frac{k\dot{x}}{j\omega} = \left(\frac{k}{j\omega}\right)\dot{x} = \left(-\frac{k}{\omega}j\right)\dot{x}$$

$$Z_{spring} = -\frac{k}{\omega}j$$

$$F = m\ddot{x} = (mj\omega)\dot{x}$$

$$Z_{mass}=mj\omega$$

$$F_{damper} = (Z_{damper})\dot{x} = C\dot{x}$$

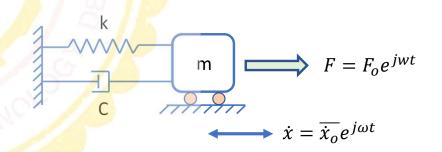
Lags
$$\dot{x}$$
 by $\frac{\pi}{2}$

$$F_{spring} = (Z_{spring})\dot{x} = -j\frac{k}{\omega}\dot{x} = (e^{-\frac{\pi}{2}j})\frac{k}{\omega}\dot{x}$$

$$pring = (Z_{spring})\dot{x} = -j\frac{k}{\omega}\dot{x} = (e^{-\frac{\pi}{2}j})\frac{k}{\omega}\dot{x}$$

$$F_{mass} = (Z_{mass})\dot{x} = m\omega j\dot{x} = \left(e^{\frac{\pi}{2}j}\right)m\omega\dot{x}$$

Leads \dot{x} by $\frac{\pi}{2}$

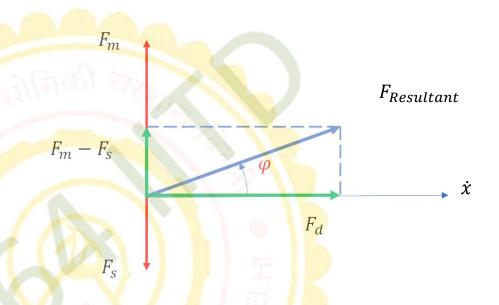




$$F_{Resultant} = \sqrt{F_d^2 + (F_m - F_s)^2}$$
$$= \sqrt{C^2 + \left(m\omega - \frac{k}{\omega}\right)^2} \dot{x}$$

$$\tan(\varphi) = \frac{F_m - F_s}{F_d}$$
$$= \frac{m\omega - \frac{k}{\omega}}{C}$$

$$Z_{eff} = Z_{resultant} = \sqrt{C^2 + \left(m\omega - \frac{k}{\omega}\right)^2}$$

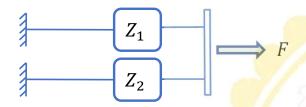


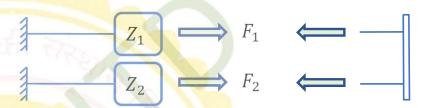
Phase Lag =
$$\varphi = \tan^{-1} \left(\frac{m\omega - \frac{k}{\omega}}{C} \right)$$

$$\dot{x} = \left(\frac{F_o}{Z_{eff}}\right) \cdot e^{j(\omega t - \varphi)}$$

Single term representation of mass, stiffness and damping

In parallel:





$$F = F_1 + F_2 \qquad \dot{x} = \dot{x}_1 = \dot{x}_2$$

$$F = F_1 + F_2 = Z_1 \dot{x} + Z_2 \dot{x} = (Z_1 + Z_2) \dot{x}$$

$$Z_{eq} = Z_1 + Z_2 + \dots + Z_n = \sum_{i=1}^{N} Z_i$$

In series

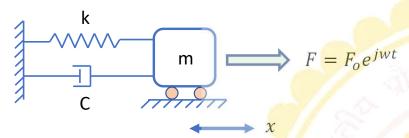
$$Z_1$$
 Z_2 F

$$F = F_1 = F_2 \qquad \dot{x} = \dot{x}_1 + \dot{x}_2$$

$$\dot{x} = \dot{x}_1 + \dot{x}_2 = \frac{F_1}{Z_1} + \frac{F_2}{Z_2} = F\left(\frac{1}{Z_1} + \frac{1}{Z_2}\right)$$

$$\frac{1}{Z_{eq}} = \frac{1}{Z_1} + \frac{1}{Z_2} + \dots + \frac{1}{Z_n} = \sum_{i=1}^{N} \frac{1}{Z_i}$$

Classical mechanics representations



 \dot{x} is same for all elements

$$F = m\ddot{x} + c\dot{x} + kx$$

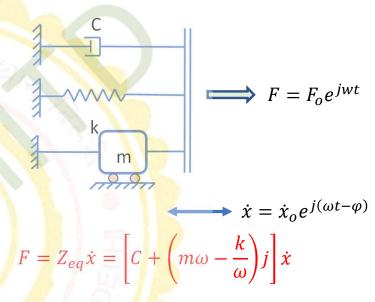
For a pure damper C(k = m = 0), $\varphi = 0$ No phase Lag

$$Z_{eq} = \sum_{i=1}^{N} Z_i = C + m\omega j - \frac{k}{\omega} j$$

$$Z_{eq} = C + \left(m\omega - \frac{k}{\omega}\right) j$$

$$Z_{eq} = X + Y j$$

Impedance mechanics

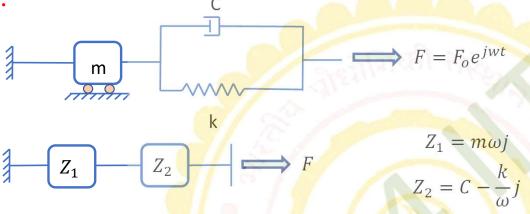


Equivalent to $(|Z|, \varphi)$ form

$$|Z| = \sqrt{C^2 + \left(m\omega - \frac{k}{\omega}\right)^2}$$

$$\varphi = \tan^{-1}\left(\frac{m\omega - \frac{k}{\omega}}{C}\right)$$





$$Z_{eq} = \frac{Z_1 Z_2}{Z_1 + Z_2} \rightarrow \text{Find out}$$

Solving differential equation eliminated. Hence computationally simpler.

Complex Electric Permittivity

Youngs modulus

$$\bar{\gamma} = \gamma_{st}(1 + \eta j)$$

$$\sigma = \sigma_o e^{j\omega t}$$

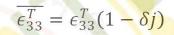
$$\epsilon = \frac{\sigma_o}{\gamma} e^{j(\omega t - \varphi)}$$

$$\sigma = \sigma_0 e^{j\omega t}$$

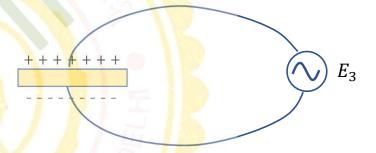
$$\epsilon = \frac{\sigma_0}{\gamma} e^{j(\omega t - \varphi)}$$



Electric permittivity



Electric Loss tangent (Electrical damping)



Electrical changes lag behind the field

$$D_{3} = \overline{\epsilon_{33}^{T}} E_{3} + d_{31} \mathcal{T}_{1}$$

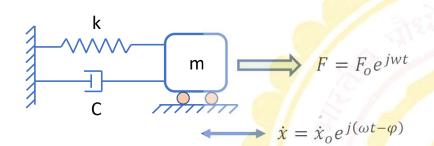
$$D_{3} = \epsilon_{33}^{T} (1 - \delta j) \cdot E_{o} e^{j\omega t}$$

$$= \epsilon_{33}^{T} E_{o} \cdot e^{j(\omega t - \delta)} \qquad e^{-j\delta} \text{ for } \delta \ll 1$$

$$= D_{3o} \cdot e^{j(\omega t - \delta)}$$

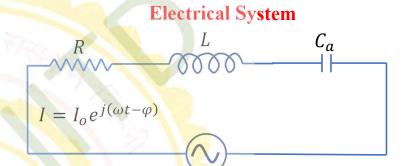
Analogy

Mechanical System



Equivalent:

$$Z = C + j \left(m\omega - \frac{k}{\omega} \right)$$



 $V = V_o e^{jwt}$

$$V \Leftrightarrow F$$
 $\dot{x} \Leftrightarrow I$

$$Z_e = R + j\left(L\omega - \frac{1}{C_a\omega}\right)$$

Equivalent:
$$C \Leftrightarrow R$$
 (Dissipative)

$$m \Leftrightarrow L$$
 (Non- Dissipative) $k \Leftrightarrow C_a$

If
$$(m,k) = 0$$
 or $(L, C_a) = 0 \implies$ No phase lag