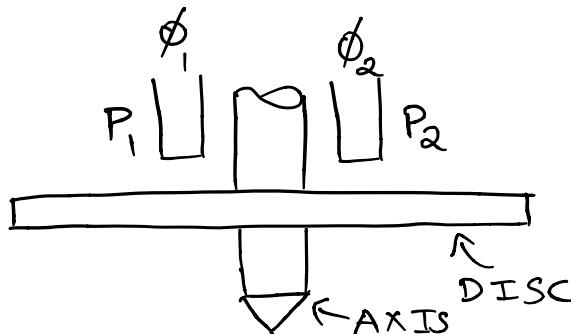


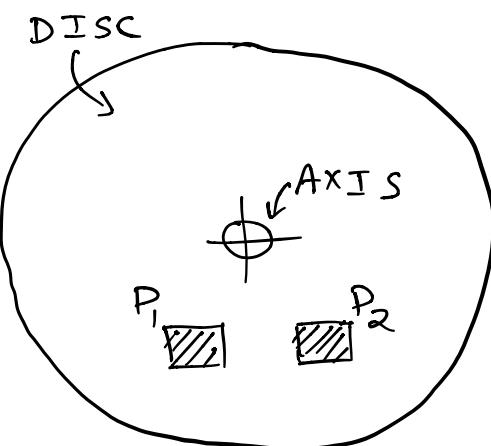
1. A thin aluminium disc is free to rotate about an axis through its centre. Two poles P_1 and P_2 produce alternating magnetic fluxes ϕ_1 and ϕ_2 cutting through the disc. If the disc is purely

$$\phi_1 = \phi_m \sin \omega t,$$

$$\phi_2 = \phi_m \sin(\omega t - \beta)$$



FRONT VIEW



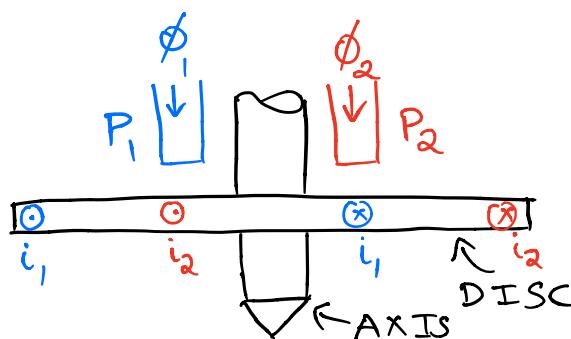
TOP VIEW

resistive, what is the torque acting on the disc?

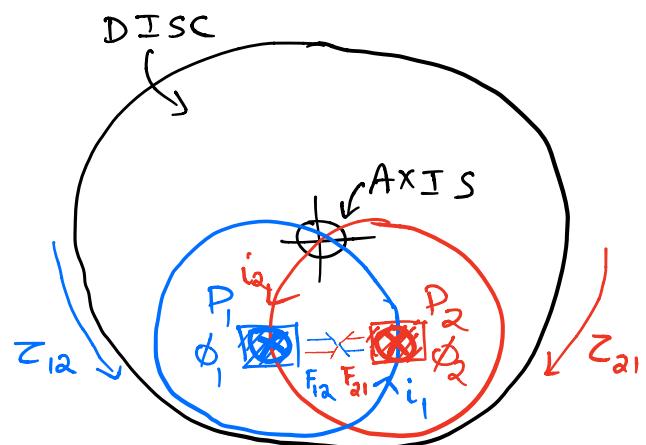
What is the torque if the disc has impedance $Z = |Z| e^{j\alpha}$? For what value of β and α is the torque maximum?

$$\phi_1 = \phi_m \sin \omega t,$$

$$\phi_2 = \phi_m \sin(\omega t - \beta)$$



FRONT VIEW



TOP VIEW

- Assuming flux directions at an instant as above, the currents induced in the disc may be obtained.

$$i_1 = \frac{\phi_m}{R} \omega \cos \omega t, i_2 = \frac{\phi_m}{R} \omega \cos(\omega t - \beta) \quad (1)$$

disc resistance

- Force act in indicated directions and produce torques.

$$F_{21} \sim \frac{\phi_m}{R} \omega \cos \omega t \times \phi_m \sin(\omega t - \beta) \rightarrow z_{21} \quad (1)$$

$$F_{12} \sim \frac{\phi_m}{R} \omega \cos(\omega t - \beta) \times \phi_m \sin \omega t \rightarrow z_{12}$$

- Net torque, $\tau \sim z_{12} - z_{21}$

$$\begin{aligned} &= \frac{\phi_m^2}{R} \omega [\cos(\omega t - \beta) \sin \omega t - \cos \omega t \sin(\omega t - \beta)] \\ &= \frac{\phi_m^2 \omega}{R} \sin \beta \end{aligned} \quad (1)$$

If disc has impedance $z = |z| e^{j\alpha}$

$$i_1 = \frac{\phi_m}{|z|} \omega \cos(\omega t - \alpha) \quad (1)$$

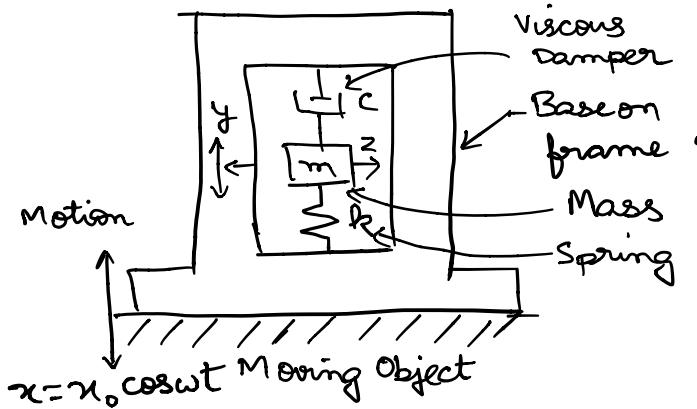
$$i_2 = \frac{\phi_m}{|z|} \omega \cos(\omega t - \beta - \alpha) \quad (1)$$

$$\Rightarrow \tau \sim \frac{\phi_m^2 \omega}{|z|} [\cos(\omega t - \beta - \alpha) \sin \omega t - \cos(\omega t - \alpha) \sin(\omega t - \beta)]$$

$$= \frac{\phi_m^2 \omega}{|z|} \sin \beta \cos \alpha \quad (1)$$

Maximum when $\beta = 90^\circ$ (fluxes at right angles),
 $\alpha = 0^\circ$ (purely resistive) (1)

2. A schematic diagram of a seismic transducer is



The equation of motion of mass 'm' is,

$$m\ddot{z} = -c(\dot{y} - \dot{z}) - k(y - z).$$

For its motion, $z = y - x$,

relative to the frame,

this equation is $m\ddot{z} + c\dot{z} + kz = -m\dot{x} = m\omega^2 x_0 \cos\omega t$.

a) Find the steady-state solution $z(t)$.

b) If Z_0 is the amplitude of $z(t)$, show that

$$\frac{Z_0}{x_0} = \frac{\gamma^2}{\sqrt{(1-\gamma^2)^2 + (2\xi\gamma)^2}}, \quad \gamma = \frac{\omega}{\omega_n}, \quad \omega_n = \sqrt{\frac{k}{m}}, \quad \xi = \frac{c}{2\sqrt{km}}$$

c) Find $\lim_{\gamma \rightarrow \infty} \frac{Z_0}{x_0}$ (for constant ξ).

d) Find $\lim_{\gamma \rightarrow 0} \frac{\omega_n^2 Z_0}{\omega^2 x_0}$ (for constant ξ).

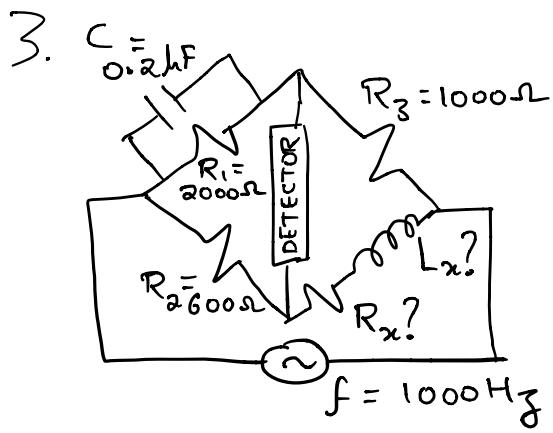
a) $z(t) = Z_0 \cos(\omega t - \phi)$

$$Z_0 = \frac{m\omega^2 x_0}{\sqrt{(k-m\omega^2)^2 + (\xi\omega)^2}}, \quad \phi = \tan^{-1}\left(\frac{\xi\omega}{k-m\omega^2}\right)$$

b) $\frac{Z_0}{x_0} = \frac{m\omega^2}{\sqrt{(k-m\omega^2)^2 + (\xi\omega)^2}} = \frac{\gamma^2}{\sqrt{(1-\gamma^2)^2 + (2\xi\gamma)^2}}$

c) $\lim_{\gamma \rightarrow \infty} \frac{Z_0}{x_0} = 1 \Rightarrow$ for large enough γ , $Z_0 \sim x_0$

d) $\lim_{\gamma \rightarrow 0} \frac{\omega_n^2}{\omega^2} \frac{Z_0}{x_0} = 1 \Rightarrow$ for small enough γ , accelerations are same



Determine the inductive impedance (R_x, L_x) connected in this balanced (Maxwell's) bridge circuit.

In balanced condition,

$$\frac{R_1 \parallel \frac{1}{j\omega C}}{R_2} = \frac{R_3}{R_x + j\omega L_x} \quad \text{①}$$

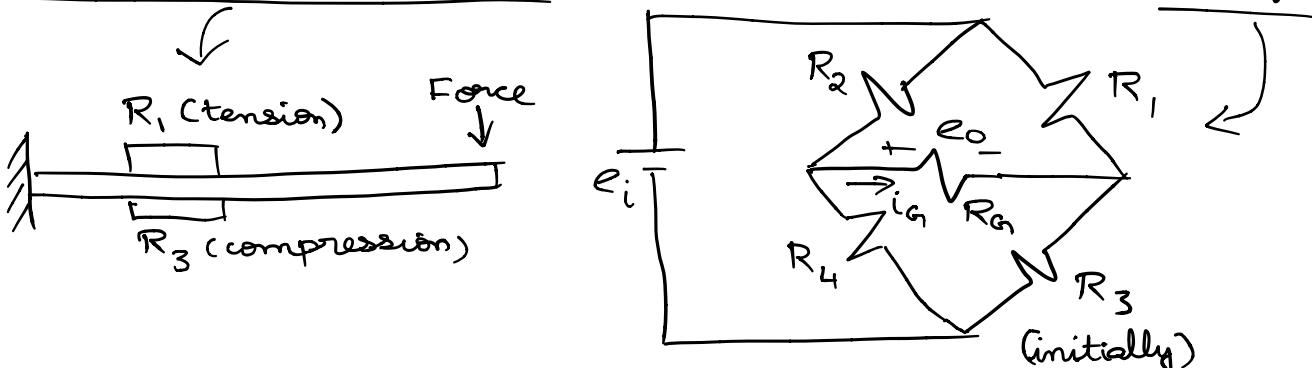
$$\Rightarrow Z_x = \frac{R_2 R_3}{R_1} + j\omega R_2 R_3 C \quad |$$

$$\Rightarrow R_x = \frac{600 \times 1000}{2000} = 300 \Omega$$

$$\delta \Rightarrow L_x = 600 \times 1000 \times 0.2 \times 10^{-6} \quad \text{①}$$

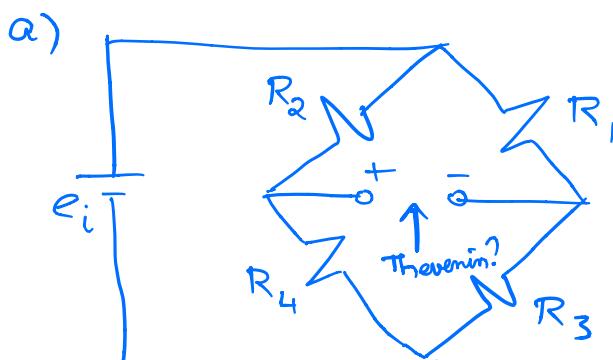
$$= 12 \text{ mH}$$

4. Two electrical strain gauges are bonded to a Duralumin cantilever and connected in a bridge

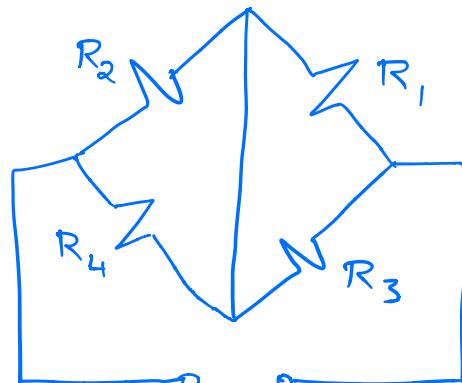


Each gauge has a resistance of 100Ω and a gauge factor of 2.1. $R_2 = R_4 = 100\Omega$. $\epsilon_i = 4V$. Stress is $200 \times 10^6 N/m^2$. Young's modulus of Duralumin 70 GN/m^2 .

- If $R_g = 10k\Omega$, $\epsilon_0 = ?$
- Comment on the temperature compensation ability of this arrangement.



$$V_{th} = \epsilon_i \left[\frac{R_4}{R_2 + R_4} - \frac{R_3}{R_1 + R_3} \right]$$



$$R_{th} = R_2 \parallel R_4 + R_1 \parallel R_3$$

$$\gamma = \frac{\text{Stress}}{\text{strain}} \Rightarrow \text{Strain} = \frac{200 \times 10^6}{70 \times 10^9} = 2.86 \times 10^{-3} \quad ①$$

$$\Rightarrow \frac{\Delta R_1}{R_1} = 2.1 \times 2.86 \times 10^{-3} = 6 \times 10^{-3} = -\frac{\Delta R_3}{R_3}$$

$$\therefore R_1 = 100(1 + 6 \times 10^{-3}), R_3 = 100(1 - 6 \times 10^{-3}) \quad ①$$

$$\Rightarrow V_{th} = 12mV, R_{th} = 99.998\Omega$$

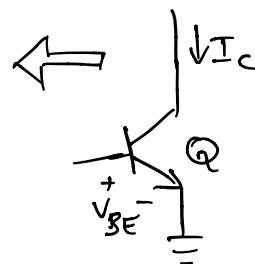
$$\therefore E_o = 12 \text{ mV} \cdot \frac{10k\Omega}{10k\Omega + 99.9982} = 11.88 \text{ mV} \quad (1)$$

(b) Arrangement is temperature compensated as the strain gauges are on adjacent arms.

(1)

5. Properties of Silicon may be used to make a temperature sensor where a voltage is Proportional To Absolute Temperature.

a) Consider a BJT,



$$\text{Assume, } I_c = I_s e^{-qV_{BE}/kT}$$

T: Absolute Temperature,

q : electron charge, k : Boltzmann constant,

$$I_s = I_0 e^{-qV_B/kT}$$

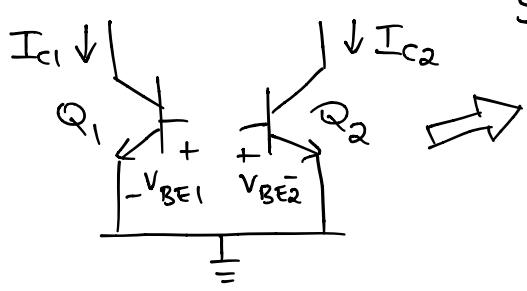
V_g : Silicon bandgap (temperature independent)

I_0 : device parameter (assumed temperature independent)

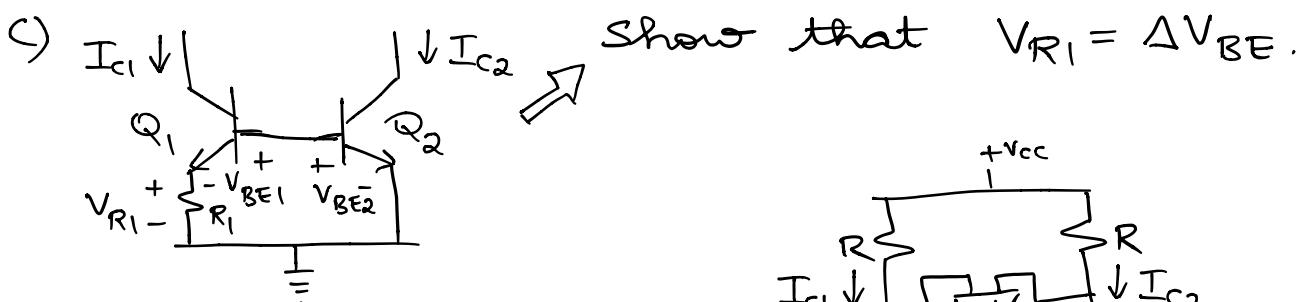
$I_0 \propto A$, area of transistor

$$\text{Show that } V_{BE} = V_g + \frac{kT}{q} \ln \frac{I_s}{I_0}, \quad \frac{\partial V_{BE}}{\partial T} = \frac{V_{BE} - V_g}{T}$$

b) Consider two BJT's with $A_1 = 8A_2$ (area relation).

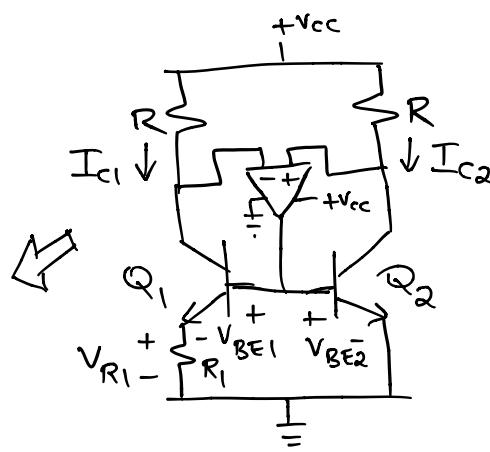


$$\begin{aligned} \text{Show that } \Delta V_{BE} &\stackrel{\text{defn.}}{=} V_{BE2} - V_{BE1} \\ &= \frac{kT}{q} \ln \left(\frac{I_{c2}}{I_{c1}} \cdot 8 \right) \end{aligned}$$



$$\text{Show that } V_{R1} = \Delta V_{BE}.$$

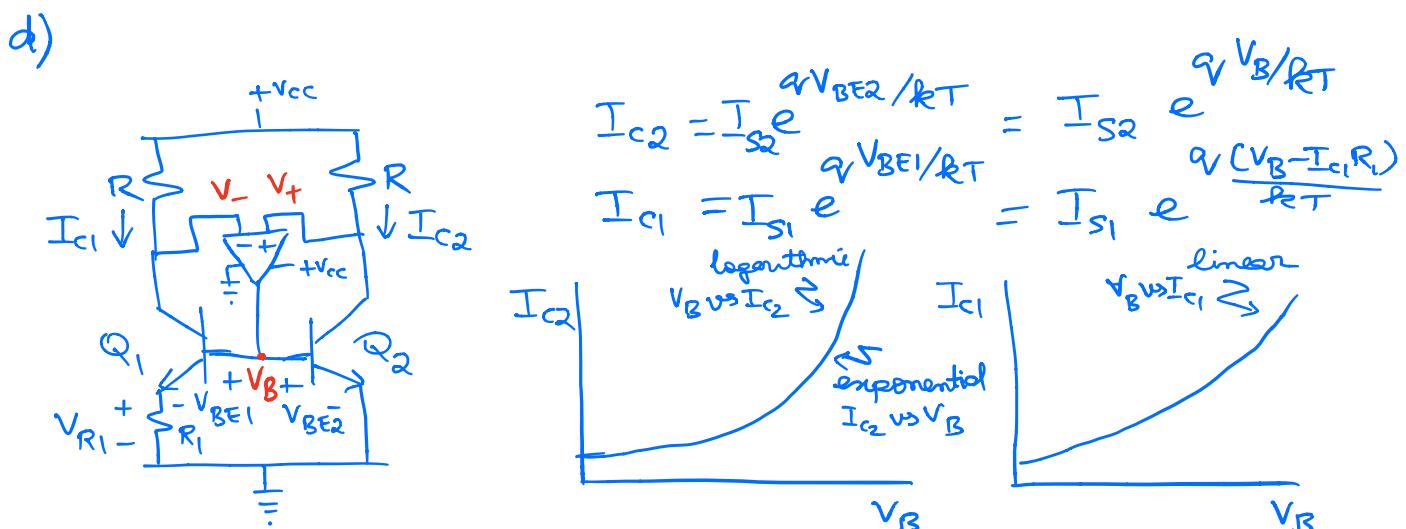
d) Discuss whether the op-amp feedback can ensure $I_{c1} = I_{c2}$.



$$\begin{aligned}
 a) \quad I_c &= I_s e^{\frac{qV_{BE}}{kT}} \\
 &= I_0 e^{-\frac{qV_B}{kT}} e^{\frac{qV_{BE}}{kT}} \\
 \Rightarrow \ln \frac{I_c}{I_0} &= \frac{V_{BE} - V_B}{kT/q} \Rightarrow V_{BE} = V_B + \frac{kT}{q} \cdot \ln \frac{I_c}{I_0} \\
 \frac{\partial V_{BE}}{\partial T} &= \frac{k}{q} \ln \frac{I_c}{I_0} = \frac{V_{BE} - V_B}{T} \quad \textcircled{1}
 \end{aligned}$$

$$\begin{aligned}
 b) \quad \Delta V_{BE} &\stackrel{\Delta}{=} V_{BE2} - V_{BE1} \\
 &= V_B + \frac{kT}{q} \ln \frac{I_{c2}}{I_{o2}} - \left(V_B + \frac{kT}{q} \ln \frac{I_{c1}}{I_{o1}} \right) \quad \textcircled{1} \\
 &= \frac{kT}{q} \ln \frac{I_{c2} \cdot I_{o1}}{I_{c1} \cdot I_{o2}} = \frac{kT}{q} \ln \frac{I_{c2}}{I_{c1}} \cdot \frac{A_1}{A_2} = \frac{kT}{q} \ln \left(\frac{I_{c2}}{I_{c1}} \cdot 8 \right)
 \end{aligned}$$

$$\begin{aligned}
 c) \quad V_{BE2} &= V_{BE1} + V_{RI} \\
 \Rightarrow V_{RI} &= V_{BE2} - V_{BE1} = \Delta V_{BE} \quad \textcircled{1}
 \end{aligned}$$



If V_B is very small, I_{c1}, I_{c2} are small, $\frac{I_{c1}}{I_{c2}} = \frac{I_{s1}}{I_{s2}} \cdot e^{-\frac{q(I_{c1}R_i + V_B)}{kT}} \approx 1$

$\Rightarrow V_- < V_+ \Rightarrow V_B$ increases

If V_B is very large, I_{c1}, I_{c2} are large and $I_{c2} > I_{c1}$

$\Rightarrow V_+ < V_- \Rightarrow V_B$ decreases

\therefore Can settle to a point where $I_{c1} = I_{c2}$

For BJT, assume $I_c = I_s e^{\frac{V_{BE}}{V_T}}$, where I_c is collector current, V_{BE} is base-emitter voltage, $V_T = \frac{kT}{q}$ is thermal voltage, and I_s is a parameter that is

process and temperature dependent $I_s = I_0 e^{-\frac{V_{ao}}{V_T}}$, I_0 is device parameter (temperature independent), V_{ao} is bandgap of silicon (1.1eV). Show that

$$V_{BE} = V_{ao} + V_T \ln \frac{I_c}{I_0} \quad \frac{\partial V_{BE}}{\partial T} = \frac{k}{q} \ln \frac{I_c}{I_0} = \frac{V_{BE} - V_{ao}}{T}$$

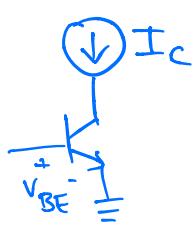
For two BJTs show that $\Delta V_{BE} = V_{BE1} - V_{BE2} = V_T \ln \frac{I_{c1}}{I_{c2}} \cdot \frac{I_{s2}}{I_{s1}}$.

Show that $V_{R1} = \Delta V_{BE}$ and $V_{R2} \propto V_T$

* Opamps.

The temperature dependence of silicon properties can be used to make temperature sensors.

Consider a BJT, Assume $I_c = I_s e^{\frac{V_{BE}}{V_T}}$,

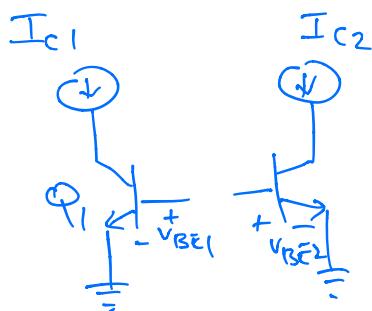


where I_c is collector current, V_{BE} is Base-Emitter voltage, $V_T = \frac{kT}{q}$ is Boltzmann's constant,

T is absolute temperature, q is charge of electron, $I_s = I_0 e^{-\frac{E_g}{kT}}$, I_0 is a device parameter (temperature independent) and E_g is bandgap of silicon? Show that

$$V_{BE} = V_{GO} + V_T \ln \frac{I_C}{I_0}, \quad \frac{\partial V_{BE}}{\partial T} = \frac{V_{BE} - V_{GO}}{T}.$$

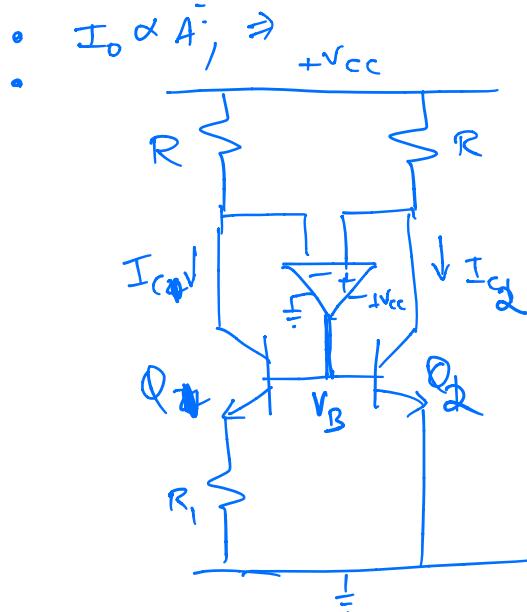
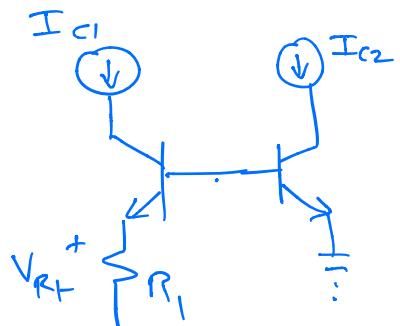
- Consider two BJTs and show that $\Delta V_{BE} = V_{BE1} - V_{BE2}$



$$= V_T \ln \frac{I_{C1}}{I_{C2}} \cdot \frac{I_{S2}}{I_{S1}}$$

$$= V_T \ln \frac{I_{S1}}{I_{S2}} \frac{I_{Q2}}{I_{Q1}}$$

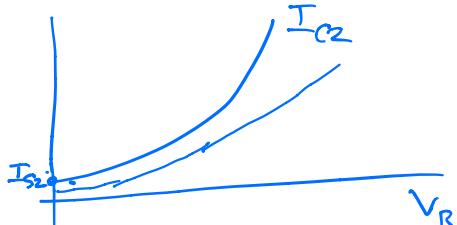
- Show that $V_{R1} = \Delta V_{BE}$



Show that the op-amp ensuring $I_{C1} = I_{C2}$
Discuss whether $A_1 = 8A_2$ if $A_1 = 8A_2$

$$I_{C1} = I_{S1} e^{\frac{V_B - I_{C1}R_1}{V_T}}$$

$$I_{C2} = I_{S2} e^{\frac{V_B - I_{C2}R_2}{V_T}} \\ = I_{S2} e^{\frac{V_B}{V_T}}$$



$$V_B = \ln \frac{I_{C1} + I_{C1}R_1}{I_{S1}}$$

$$V_B = V_T \ln \frac{I_{C1}}{I_{C2}}$$



Suppose $I_{C1} \neq I_{C2}$.

$I_{C1} > I_{C2}$ or $I_{C1} < I_{C2}$

$\Rightarrow V_- < V_+$

$\Rightarrow V_B > 0$

$I_Q \uparrow \Rightarrow V_- \downarrow \Rightarrow V_B \uparrow \Rightarrow I_{C2} \uparrow$

$I_{C1} \downarrow \Rightarrow V_- \uparrow \Rightarrow V_B \downarrow \Rightarrow I_{C2} \downarrow$

$I_{C2} \uparrow \Rightarrow V_+ \downarrow \Rightarrow V_B \downarrow \Rightarrow \frac{I_{C2}}{I_{C1}} \downarrow$