# Math 2E03- Introduction to Modelling 

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## The Wave Equation

Consider, if you will, a violin string of length $L$ and linear density (mass per length) $\sigma$, under a tension $T$. If the string runs along the $x$ axis, we may define the displacemnet of the string as it vibrates as $U(x, t)$ as shown in the graph (given in class). If $\Delta x$ is small enough, we may assume that the angles are quite small and that the horizontal componnets of the forces cancel each other out. This leaves the vertical forces $T \sin \theta_{2}-T \sin \theta_{1}$. These forces must equal to given $m a$, where the mass is given by the linear density times the length and the acceleration may be written in differential form:

$$
T \sin \theta_{2}-T \sin \theta_{1} \approx \sigma \Delta x \frac{\partial^{2} U}{\partial t^{2}}
$$

For very small angles, $\sin x \approx x \approx \tan x$ and in this case we will use the tangents since they may always be written as the difference of hight $\partial U$ divided by the difference in width $\partial x$ :

$$
\begin{aligned}
& T \tan \theta_{2}-T \tan \theta_{1} \approx \sigma \Delta x \frac{\partial^{2} U}{\partial t^{2}} \\
& \frac{\frac{\partial U(x+\Delta x)}{\partial x}-\frac{\partial U(x)}{\partial x}}{\Delta x} \approx \frac{\sigma}{T} \frac{\partial^{2} U}{\partial t^{2}} \\
& \frac{\partial^{2} U}{\partial t^{2}}=c^{2} \frac{\partial^{2} U}{\partial x^{2}} \text { where } c=\sqrt{\frac{T}{\sigma}}
\end{aligned}
$$

Solving Wave equation using Fourier's method

Again here we are looking for a solution of the form $\tilde{U}(x, t)=X(x) T(t)$

$$
\text { In short form } \begin{aligned}
U & =X T \\
U_{t} & =X \dot{T} \\
U_{t t} & =X \ddot{T} \\
U_{x} & =\dot{X} T \\
U_{x x} & =\ddot{X} T
\end{aligned}
$$

Substituting these values of $U_{t t}$ and $U_{x x}$ in wave equation we get

$$
\begin{gathered}
\Longrightarrow X \ddot{T}=c^{2} \ddot{X} T \\
\\
\frac{\ddot{T}}{T}=c^{2} \frac{\ddot{X}}{X}
\end{gathered}
$$

We again have two indpendent ODEs. Now we will set upthese ODEs equal to the constant $-\lambda^{2}$.

$$
\begin{aligned}
\ddot{T}+c^{2} \lambda^{2} T & =0 \\
\ddot{X}+\lambda^{2} X & =0
\end{aligned}
$$

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So the solution of these two ODEs is

$$
\begin{aligned}
& T=A \sin (\lambda c t)+B \cos (\lambda c t) \\
& X=C \sin (\lambda x)+D \cos (\lambda x)
\end{aligned}
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In our example of violin string, our ends are fixed at all times and we will call the initial shape of our string $f(x)$. The another boundary condition is at $t=0$ velocity at all points will be zero. -Now we will add boundary values $U(0, t)=U(L, t)=0$. Let's us impose our first boundary value

$$
U(0, t)=X(0) T(t)=0
$$

$$
\Longrightarrow X(0)=0
$$

$$
\begin{aligned}
& B \sin (0)+C \cos (0)=0 \\
& B(0)+C(1)=0 \\
& C=0 \\
& \quad \Longrightarrow X=B \sin (\lambda x)
\end{aligned}
$$

We now impose the other boundary value:

$$
\begin{aligned}
U(L, t) & =X(L) T(t)=0 \\
& \Longrightarrow X(L)=0 \\
C & \sin (\lambda L)=0 \\
\lambda L & =n \pi \\
\lambda & =\frac{n \pi}{L} \\
& \Longrightarrow X_{n}=C_{0} \sin \left(\frac{n \pi}{L_{0}^{\circ}} x\right) .
\end{aligned}
$$

$$
\Longrightarrow T_{n}=A_{n} \sin (\lambda c t)+B_{n} \cos (\lambda c t)
$$

$$
\Longrightarrow U_{n}=\left[A_{n} \sin \left(\frac{n \pi}{L} c t\right)+B_{n} \cos \left(\frac{n \pi}{L} c t\right)\right] \sin \frac{n \pi}{L} x
$$

