

Hydrogen atom

$$\hat{H} = -\frac{\hbar^2}{2\mu} \nabla^2 - \frac{Ze}{r}$$

Spherical coordinates

$$\nabla^2 = \frac{1}{r^2} \frac{\partial}{\partial r} r^2 \frac{\partial}{\partial r} + \frac{1}{r^2 \sin \theta} \frac{\partial \sin \theta}{\partial \theta} \frac{\partial}{\partial \theta} + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2}{\partial \phi^2}$$

L²

$$\Psi(r, \theta, \varphi) = R(r) Y_{lm}(\theta, \varphi)$$

$$L^2 Y_{lm} = l(l+1) \hbar^2 Y_{lm}$$

$$\hat{H} = \left(-\frac{\hbar^2}{2\mu r^2} \frac{\partial}{\partial r} r^2 \frac{\partial}{\partial r} + \frac{l(l+1)\hbar^2}{2\mu r^2} - \frac{Ze^2}{r} \right)$$

↑ centrifugal

$$\hat{H} R = E R$$

$$\text{For } l \neq 0, \quad V_{\text{eff}}(r) = \frac{l(l+1)}{r^2} - \frac{Ze^2}{r}$$

$$V_{\text{eff}}(r) \rightarrow +\infty \quad \text{as } r \rightarrow 0$$

$$R(r) \rightarrow 0 \quad \text{as } r \rightarrow 0$$

$$\text{For } l = 0, \quad V_{\text{eff}}(r) = -\frac{Ze^2}{r} \rightarrow -\infty \text{ as } r \rightarrow 0$$

$$R(r) = \underline{g(r)}$$

$$r = \frac{\hbar^2}{\mu Ze^2} \rho \rightarrow a_0 \text{ Bohr radius}$$

$$\bar{E} = \frac{Ze^2}{a_0} \epsilon \rightarrow \text{atomic unit of energy}$$

$\rho \rightarrow$ dimensionless length

$\epsilon \rightarrow$ dimensionless energy

$$-\frac{d^2 g}{dp^2} + \frac{l(l+1)g}{p^2} - \frac{2}{p}g = 2\epsilon g$$

Case A: For $l=0$

Step 1: Look for solution at large p
 $g(p) \sim e^{-bp}$ $b^2 = -2\epsilon$

$$g(p) = p e^{-bp} \text{ so that } g(p) = 0 \text{ at } p=0$$

$$R(r) = \frac{g(r)}{r}$$

$$R(p) = \frac{g(p)}{p}$$

No nodes $g_1(\rho) : \rho e^{-b\rho}$

$g_1(\rho)$ is an eigenfunction of the radial Schrödinger equation when

$$2E_1 = -1 \quad \text{or} \quad \boxed{E_1 = -\frac{1}{2}} \quad l=0$$

$\frac{1}{n^2}$

one node

$$g_2(\rho) = (c_0\rho - c_1\rho^2) e^{-b\rho}$$

$$E_2 = -\frac{1}{4} = \frac{1}{4} E_1$$

two nodes

$$g_3(\rho) = (a\rho + b\rho^2 + c\rho^3) e^{-b\rho}$$

$$E_3 = \frac{1}{9} E_1$$

Case B: $l \neq 0$

for large ρ $g(\rho) \sim e^{-b\rho}$

for small ρ , $g(\rho) \sim \rho^{-l}$ or ρ^{l+1}

$$g(\rho) = \rho^{l+1} e^{-b\rho} \times \text{polynomial in } \rho$$

$$E = -\frac{1}{2} \frac{1}{(l+1)^2}$$

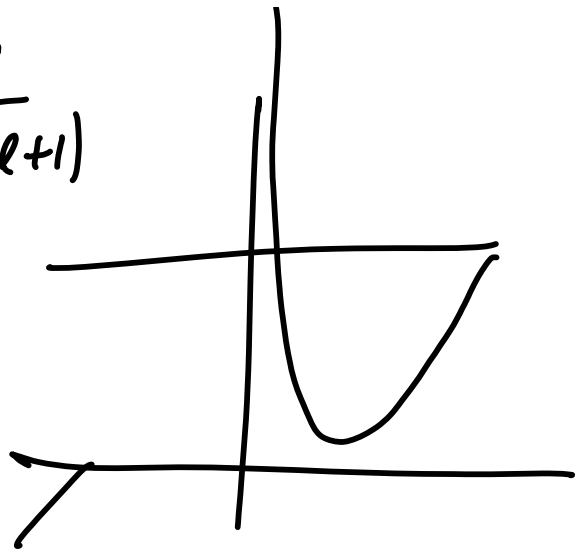
$$V_{\text{eff}} = \frac{l(l+1)}{r^2} - \frac{ze^2}{r}$$

$$V_{\text{eff}}^{\text{min}} = \frac{1}{\ell(\ell+1)}$$

$$\ell < 0 \rightarrow n-1$$

$$\frac{dV_{\text{eff}}}{d\rho} = 0$$

solve for ρ



$$\Psi_{n\ell m} = R_{n\ell}(r) Y_{\ell m}(\theta, \phi)$$

$R_{n\ell}(r) = r^{\ell} e^{-\frac{Zr}{na_0}} \times \text{polynomial in } r$

ℓ → orbital angular momentum #
 m → no. of nodes
 (n, ℓ, m) → principal quantum number

$$R(r) = e^{-Zr/a_0} \times 1$$

$$\int R_{n\ell}^*(r) R_{n\ell}(r) r^2 dr = 1$$

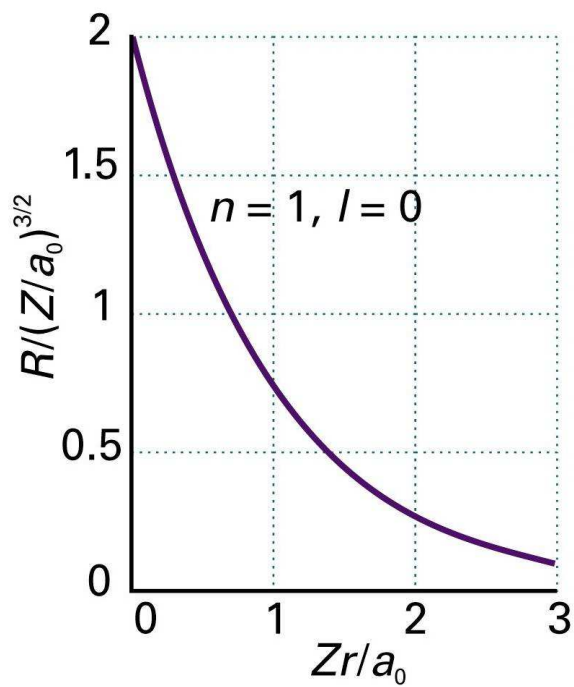


Figure 10-4a
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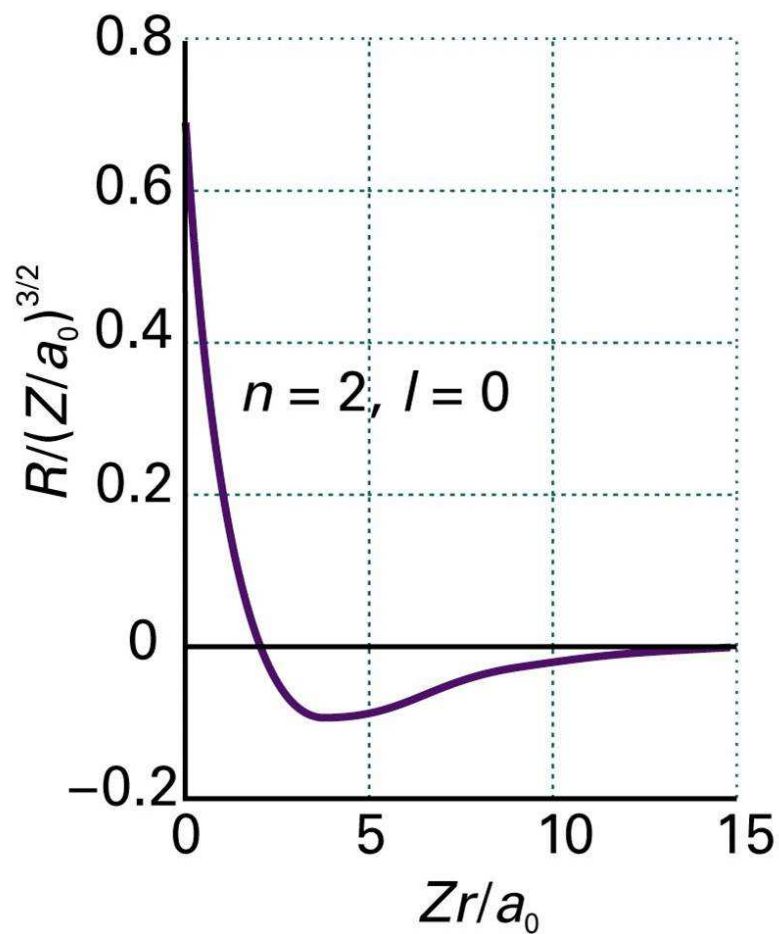


Figure 10-4b
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$$R_{21} = e^{-\frac{Zr}{2a_0}} \times \rho$$

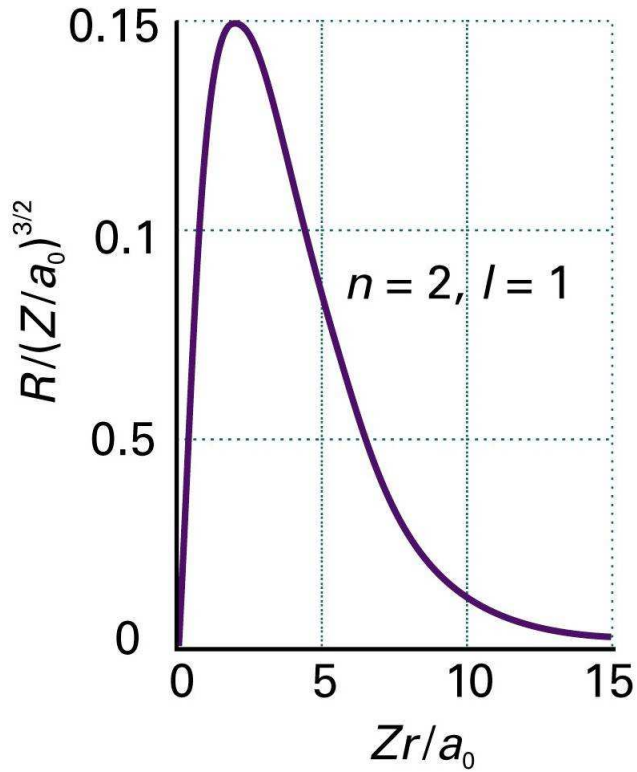


Figure 10-4d
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$$\langle r \rangle = \int R_{nl}^*(r) r R_{nl}(r) r^2 dr \int Y_{lm}^* Y_{lm} \sin\theta d\theta d\phi$$

$$\langle O_p \rangle = \int \psi^* O_p \psi dz$$

$$\langle r \rangle = \int R_{nl}^*(r) Y_{lm}^*(\theta, \phi) r Y_{lm}(\theta, \phi) r^2 dr d\theta d\phi$$

$$= \underbrace{\int R_{nl}^*(r) R_{nl}(r) r^2 dr}_1 \times \underbrace{\int \int Y_{lm}^*(\theta, \phi) Y_{lm}(\theta, \phi) \sin\theta d\theta d\phi}_1$$

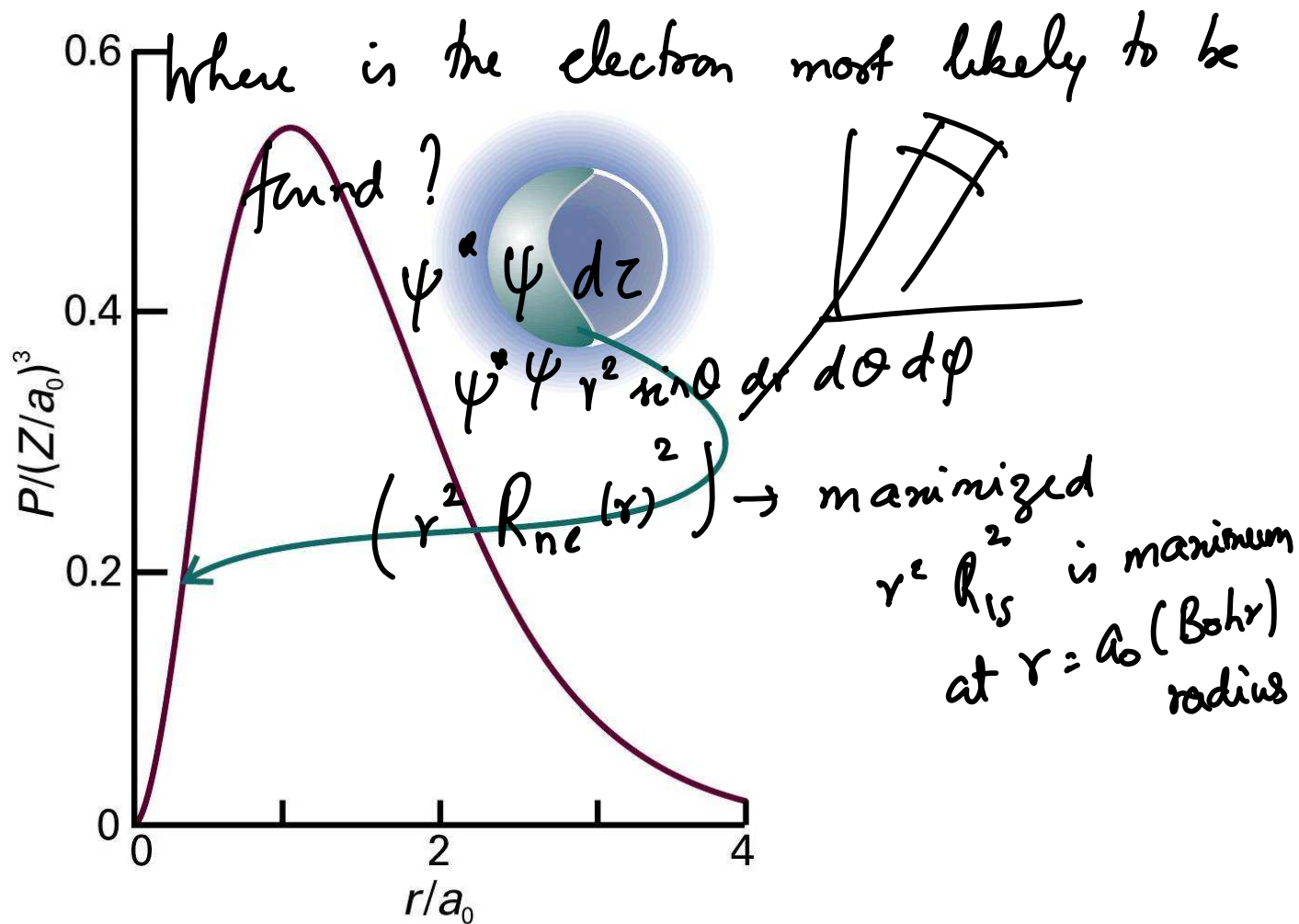


Figure 10-14
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$$\Psi = R(r) Y_{lm}(\theta, \varphi)$$

$$\Psi^* \Psi dr = R^2 |Y_{lm}|^2 r^2 dr \sin\theta d\theta d\varphi$$

$$\langle r \rangle \neq \frac{V_{mp}}{\sqrt{\int R_{ne}(r) r^2 R_{ne}(r) dr}}$$

Statistics of the miner

Class average = 32

Std dev = 14

Max = 67

Min < 0