

discrete set of frequencies which contribute to the total power have changed from  $n/2\pi$  c.p.s. to  $n/2T$  c.p.s.,  $n = 0, 1, 2, \dots$ ; as shown in Fig. 4.3.

Note that if  $T$  is much larger than  $\pi$  the set of frequencies ( $1/2T, 2/2T, 3/2T, \dots$ ), which form the above power spectrum is much more "closely packed" than the standard set, ( $1/2\pi, 2/2\pi, 3/2\pi, \dots$ ), which form the power spectrum for functions with periodicity  $2\pi$ .

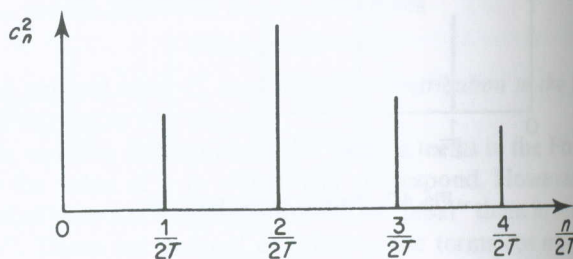


Fig. 4.3. Discrete power spectrum.

#### 4.5 NON-PERIODIC FUNCTIONS; FOURIER INTEGRALS

We now turn our attention to the case of a real valued non-periodic function  $X(t)$  (but still deterministic). Since  $X(t)$  does not now possess any form of periodic structure we cannot express it in the form of a Fourier series which would be valid for *all*  $t$ . However, we can construct a Fourier series which will represent it over a *finite* interval by using the standard trick of defining a new function which is identical to  $X(t)$  over a certain interval, but is periodic outside this interval. We thus choose an interval,  $(-T, T)$ , and define a new function,  $X_T^*(t)$ , by

$$\begin{aligned} X_T^*(t) &= X(t), & -T \leq t \leq T, \\ X_T^*(t + 2pT) &= X_T^*(t), & p = 1, 2, 3, \dots \end{aligned} \quad (4.5.1)$$

Then  $X_T^*(t)$  is certainly periodic, with period  $2T$ , and, subject to the usual regularity conditions, may be expressed in a Fourier series of the form (4.3.1), i.e. we may write,

$$X_T^*(t) = \sum_{n=0}^{\infty} (a_n \cos 2\pi f_n t + b_n \sin 2\pi f_n t), \quad (4.5.2)$$

where, for each  $n$ ,  $f_n = n/2T$ , and  $b_0 = 0$ .