

EEL-306.

(1)

LECTURES- 28 & 29.

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Response of the PLL to a sinusoidal frequency deviation from the carrier frequency, and study of second order PLLs.

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In the previous lecture we had studied the response of a first order PLL to an input which had a ~~φ~~ constant frequency shift from the carrier frequency  $f_c$  (corresponds to  $m(t) = \text{step fn}$ ). We had for  $t \gg 2T_p$ ,

$$-\frac{d\phi_e(t)}{dt} = a \sin \phi_e(t) + \pi k_{vco} m(t - \tau), \text{ where}$$

$$a \triangleq \pi k_{vco} A_c A_v, \text{ and we}$$

also had the initial condition  $\phi_e(t = 2T_p) = 0$ .

In this lecture we will extend our analysis to the case where

$$m(t) = \begin{cases} A_m \cos 2\pi f_m (t - T_p), & t \geq T_p \\ 0, & t < T_p \end{cases}$$

With the assumption that

$\phi_e(t)$  is small, we have

$\sin \phi_e(t) \approx \phi_e(t)$  — (2), using this in (1)

we set  $-\frac{d\phi_e(t)}{dt} = a\phi_e(t) + m_{kf} m(t-z)$ ,  $t > z + T_p$

with  $\phi_e(t = z + T_p) = 0$ . — (3)

Note that (3) is a linear differential equation, which can be solved.

The solution to (3) gives us a  $\phi_e(t)$  which is small for all  $t > z + T_p$ ,

then we can argue that our approximate analysis is correct (based on the fact that  $\sin(\phi_e(t)) \approx \phi_e(t)$  for small  $\phi_e(t)$ ).

We next try to solve (3).

Taking the Fourier transform of both sides of (3) we get

~~$-jmf \phi_e(f) = a\phi_e(f) + bM(f)e^{jmfz}$~~

~~$\therefore \phi_e(f) = \frac{-M(f)e^{jmfz}}{(a + jmf)}$~~

3  
 Taking since we are only interested  
 in ~~for~~  $t \geq 2+T_p$ , we assume that

$\phi_e(t) = 0$  for  $t \leq 2+T_p$ . Now, taking the  
 Fourier transform of both sides of (3) we

get

$$-j2\pi f \phi_e(f) = a \phi_e(f) + 2\pi k_f M(f) e^{-j2\pi f 2}$$

$$\text{i.e., } \phi_e(f) = \frac{-2\pi k_f M(f) e^{-j2\pi f 2}}{(a + j2\pi f)} \quad \text{--- (4)}$$

Also, since in the first order PLL

$$V(t) = -\frac{A_c A_v}{2} \sin \phi_e(t)$$

$$\approx -\frac{A_c A_v}{2} \phi_e(t), \text{ we have}$$

$$V(f) \approx -\frac{A_c A_v}{2} \phi_e(f)$$

$$= -\frac{A_c A_v M(f)}{2} \left( \frac{\phi_e(f)}{M(f)} \right) \quad \text{Using (4) we get}$$

$$= \frac{A_c A_v M(f) 2\pi k_f e^{-j2\pi f 2}}{2 (a + j2\pi f)} \quad \text{--- (5)}$$



(4)

To find  $\phi_e(t)$ , we now take the inverse Fourier transform of (4).

Since  $H(f) = \frac{1}{(a + j\pi f)}$  corresponds to an LTI system with impulse response  $h(t) = e^{-at} u(t)$ ,

from (4) we can easily say that

$$\phi_e(t) = -m k_f (m(t-\tau) \otimes e^{-at} u(t))$$

↑  
convolution.

$$= -m k_f \int_0^{t-\tau-T_p} e^{-ax} m(t-\tau-x) dx$$

( $t \geq \tau + T_p$ )

This solution to (3) clearly satisfies the

initial condition

$$\phi_e(t = \tau + T_p) = 0.$$



$$\therefore \phi_e(t) = -mk_f \int_0^{t-z-T_p} e^{-ax} m(t-z-x) dx$$

which clearly satisfies the initial condition  $t \geq z+T_p$ .

$$\phi_e(t = T_p+z) \equiv 0.$$

With  $m(t) = \begin{cases} A_m \cos mfm(t-T_p) & t \geq T_p \\ 0 & \text{for } t < T_p \end{cases}$

We get

$$\begin{aligned} \therefore \phi_e(t) &= -mk_f A_m \int_0^{t-z-T_p} e^{-ax} \cos mfm(t-z-x-T_p) dx \\ &= +mk_f A_m \int_{t-z-T_p}^0 e^{a(v-(t-z-T_p))} \cos mfm v dv \end{aligned}$$

$v \equiv t-z-T_p - x$

$$\phi_e(t) = -mk_f A_m e^{-a(t-z-T_p)} \int_0^{t-z-T_p} e^{av} \cos mfm v dv.$$

$t \geq z+T_p$

using the fact that

$$\int e^{\alpha x} \cos \beta x dx = \frac{1}{\alpha^2 + \beta^2} e^{\alpha x} [\alpha \cos \beta x + \beta \sin \beta x]$$

$$\int_0^{t-z-T_p} e^{av} \cos mfm v dv$$

$$= \frac{1}{[a^2 + (mfm)^2]} \left\{ e^{a(t-z-T_p)} \left( a \cos mfm(t-z-T_p) + mfm \sin mfm(t-z-T_p) - a \right) \right\}$$

$$\therefore \phi_e(t) = \frac{-mkyAm}{[a^2 + (mfm)^2]} \left\{ a \cos mfm(t-z-T_p) + mfm \sin mfm(t-z-T_p) - a e^{-a(t-z-T_p)} \right\} \quad \text{--- (10)}$$

For all  $t \geq z+T_p$ .

$$|\phi_e(t)| = \frac{mkyAm}{a} \left| \frac{a^2 \cos mfm(t-z-T_p)}{(a^2 + (mfm)^2)} + \frac{a \cdot mfm \sin mfm(t-z-T_p)}{(a^2 + (mfm)^2)} - \frac{a^2 e^{-a(t-z-T_p)}}{(a^2 + (mfm)^2)} \right|$$

For any 2 real nos  $x, y$ ,  $\frac{x^2}{x^2+y^2} < 1$  and  $\frac{|xy|}{x^2+y^2} < \frac{1}{2}$ .

$$\therefore |\phi_e(t)| < \frac{mkyAm}{a} \left( \frac{5}{2} \right)$$

Hence, a sufficient condition which guarantees that  $|\phi_e(t)| \ll 1$  is that

$$a \gg mkyAm \quad \text{; i.e.,}$$

$$\text{i.e.,} \quad \frac{k v_{10} A_2 A_1}{2} \gg kyAm$$

for  $a \gg \omega k_f A_m$

we know that  $\phi_e(t)$  is small,

$$\therefore v(t) \approx \frac{-A_c A_v k_f A_m}{2} \phi_e(t)$$

$$= \frac{A_c A_v \omega k_f A_m}{[a^2 + (\omega k_f A_m)^2]} \left\{ \begin{aligned} &a \cos \omega k_f A_m (t - \tau - T_p) \\ &+ \omega k_f A_m \sin \omega k_f A_m (t - \tau - T_p) \\ &- a e^{-a(t - \tau - T_p)} \end{aligned} \right\} \quad \text{--- (13)}$$

In the steady state, i.e.,

when  $(t - \tau - T_p) \gg 1/a$  --- (14)

$$v(t) \approx \frac{\omega A_c A_v k_f A_m}{\sqrt{a^2 + (\omega k_f A_m)^2}} \cos \omega k_f A_m (t - \tau - T_p - \alpha)$$

where  $\alpha = \tan^{-1} \left( \frac{\omega k_f A_m}{a} \right)$  --- (15)

$\therefore$  Compared to  $m(t) = A_m \cos \omega k_f A_m (t - \tau - T_p)$ ,

$v(t)$  has an amplitude gain of

$$G(\omega) = \frac{\omega A_c A_v k_f}{\sqrt{a^2 + (\omega k_f A_m)^2}} \quad \text{and a}$$

phase lag of

$$\alpha = \tan^{-1} \left( \frac{\omega k_f A_m}{a} \right) \quad \text{--- (16)}$$



Note that the gain factor can be written as

$$G(f_m) = \frac{k_f}{k_{vco}} \frac{a}{\sqrt{a^2 + (mf_m)^2}}$$

For minimum distortion we would like

$$G(f) = \frac{k_f}{k_{vco}} \cdot \frac{a}{\sqrt{a^2 + mf^2}} \text{ to be almost}$$

flat in the band where  $m(t)$  has

significant power; and also that

the phase response

$$\alpha(f) = \tan^{-1} \left( \frac{mf}{a} \right) \text{ be almost}$$

linear w.r.t  $f$  in the band occupied

by  $m(t)$  to which  $m(t)$  is limited to.

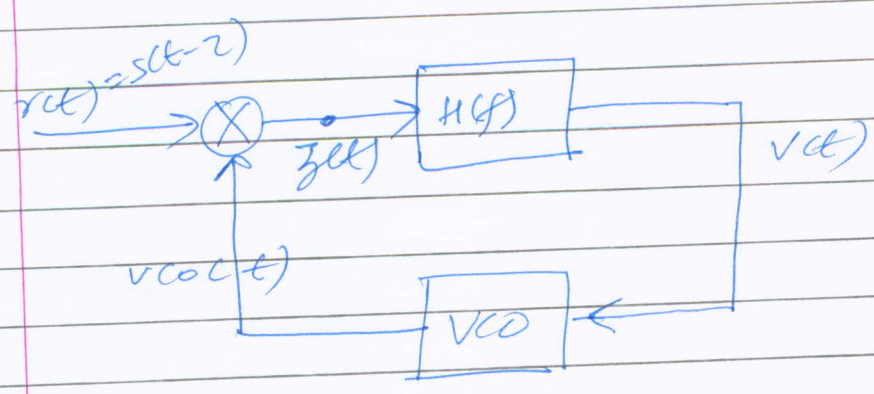
these two requirements are met simultaneously

if and only if  $\frac{a}{m} \gg$  highest frequency component of  $m(t)$

therefore we see that the same control parameter "a", controls both the bandwidth

of  $m(t)$  (such that PLL reproduces it with minimal distortion) and also controls the maximum possible freq deviation ( $k_f A_m$ ) such that  $P_e(t)$  is small. (i.e., PLL is locked)

SECOND ORDER PHASE-LOCKED LOOP.



Compared to a first order system

where  $H(f) = 1$ , in a second order system we have  $H(f) = 1 + \frac{c}{jf}$  where  $c$  is a constant. (18)

As before,

$$v_c(t) = -A_v \sin\left(2\pi f_c(t - \tau) + 2\pi k_{vco} \int v(t) dt\right)$$

$$= -A_v \sin\left(2\pi f_c(t - \tau) + \phi_2(t)\right) \quad (19)$$

and

$$r(t) = s(t - \tau)$$

$$\approx A_c \cos\left(2\pi f_c(t - \tau) + 2\pi k_f \int m(x - \tau) dx\right)$$

$$\approx A_c \cos\left(2\pi f_c(t - \tau) + \phi_1(t)\right) \quad (19)$$

the input to the low pass filter is

$$z(t) = \frac{-A_c A_v}{2} \sin(\phi_2(t) - \phi_1(t))$$

$$= \frac{-A_c A_v}{2} \sin\left(4\pi f_c(t - \tau) + \phi_1(t) + \phi_2(t)\right)$$

↑  
This term gets rejected by the low pass filter.



∴ the output of the LPF therefore

$$v(t) = \frac{-A_c A_v}{2} [\sin \phi_0(t)] * h(t) \quad \text{--- (21)}$$

↑  
denotes convolution

$h(t)$  is the impulse response of the LPF (H(f))

$$h(t) = \int_{-\infty}^{\infty} H(f) e^{j2\pi f t} df$$

$$\phi_e(t) = \phi_2(t) - \phi_1(t)$$

$$\therefore \frac{-d\phi_e(t)}{dt} = 2\pi k_f m(t - \tau) + \pi k_{vco} A_c A_v (\sin \phi_0(t) * h(t)) \quad \text{--- (22)}$$

Assuming that the parameters  $k_{vco}$  and the filter H(f) are such that  $\phi_0(t)$  is small then

$$\sin \phi_0(t) \approx \phi_0(t)$$

with this approximation in (22) we get

$$\frac{-d\phi_e(t)}{dt} = 2\pi k_f m(t - \tau) + \pi k_{vco} A_c A_v (\phi_e(t) * h(t)) \quad (t \gg T_p + \tau)$$

with the initial condition  $\phi_e(t) = 0$  at  $t = T_p + \tau$ . --- (23)

Taking Fourier transforms on both sides we get

$$-j2\pi f \phi_e(f) = 2\pi k_f M(f) e^{-j2\pi f \tau} + \pi k_{vco} A_c A_v \phi_e(f) H(f)$$



$$\frac{\phi_e(f)}{M(f)} = \frac{-mky e^{-j\omega f z}}{(j\omega f + nk_{vco} A_c A_v H(f))}$$

$$\left| \frac{\phi_e(f)}{M(f)} \right| = \frac{mky}{|j\omega f + nk_{vco} A_c A_v H(f)|}$$

and

$$V(f) \approx \frac{-A_c A_v}{2} (\phi_e(f) \otimes h(f)) \quad \left( \begin{array}{l} \sin(\phi_e(f)) \\ \approx \phi_e(f) \end{array} \right)$$

$$\therefore V(f) = \frac{-A_c A_v}{2} \phi_e(f) H(f)$$

$$\begin{aligned} \therefore \frac{V(f)}{M(f)} &= \frac{-A_c A_v H(f) \phi_e(f)}{2 M(f)} \\ &= \frac{nk_y A_c A_v H(f) e^{j\omega f z}}{(j\omega f + nk_{vco} A_c A_v H(f))} \end{aligned}$$

with  $H(f) = 1 + \frac{C}{jf}$ , from (24) we have:

$$\begin{aligned} \frac{\phi_e(f)}{M(f)} &= \frac{-mky e^{-j\omega f z}}{(j\omega f + nk_{vco} A_c A_v (1 + \frac{C}{jf}))} \\ &= \frac{(-mky e^{-j\omega f z}) jf}{(-\omega f^2 + nk_{vco} A_c A_v C + j nk_{vco} A_c A_v)} \end{aligned}$$

$$(-mf^2 + n k_{vco} A_2 A_V C + j n k_{vco} A_2 A_V f)$$

$$= (a_1 + jmf) (b_1 + jf)$$

$$a_1 b_1 = n k_{vco} A_2 A_V C$$

$$a_1 + 2mf = n k_{vco} A_2 A_V$$

i.e.,  $a_1 + 2m \cdot \frac{n k_{vco} A_2 A_V C}{a_1} = n k_{vco} A_2 A_V$

i.e.,  $a_1^2 - (n k_{vco} A_2 A_V) a_1 + 2m^2 k_{vco} A_2 A_V C = 0$

$$a_1 = \frac{n k_{vco} A_2 A_V \pm \sqrt{n^2 k_{vco}^2 A_2^2 A_V^2 - 8m^2 k_{vco} A_2 A_V C}}{2}$$

$$= \frac{n k_{vco} A_2 A_V \pm n \sqrt{k_{vco} A_2 A_V (k_{vco} A_2 A_V - 8C)}}{2}$$

assuming  $(k_{vco} A_2 A_V > 8C)$  we set  $a_1$  to be real valued and.

$$a_1 = \frac{n}{2} (k_{vco} A_2 A_V + \sqrt{k_{vco} A_2 A_V (k_{vco} A_2 A_V - 8C)})$$

$$\therefore b_1 = \frac{n k_{vco} A_2 A_V C}{a_1}$$

$$= \frac{2 k_{vco} A_2 A_V C}{k_{vco} A_2 A_V + \sqrt{k_{vco} A_2 A_V (k_{vco} A_2 A_V - 8C)}}$$

$$= \frac{2C}{1 + \sqrt{1 - \frac{8C}{k_{vco} A_2 A_V}}}$$



using the factorization on the previous page, for  $\text{Re } s > \sigma_c$ , we get  
∴ In (26) we have

$$\frac{P(s)}{M(s)} = \frac{(-ank_f) e^{-jmfz}}{(a_1 + jmf)(b_1 + jf)}$$

where  $a_1, b_1$  are given in (27).

$$\begin{aligned} \therefore \frac{-P(s)}{M(s)nk_f e^{-jmfz}} &= \frac{jf}{(a_1 + jmf)(b_1 + jf)} \\ &= \frac{d_1}{(a_1 + jmf)} + \frac{d_2}{(b_1 + jf)} \end{aligned}$$

$$d_1 b_1 + d_2 a_1 = 0 \text{ and}$$

$$d_1 + m d_2 = 1$$

$$\therefore d_1 + m \left( -\frac{d_1 b_1}{a_1} \right) = 1$$

$$\therefore d_1 = \frac{1}{(1 - mb_1/a_1)} \text{ and}$$

$$d_2 = \frac{-b_1/a_1}{(1 - mb_1/a_1)} \quad \text{--- (28)}$$

$$\therefore \frac{-P(s)}{nk_f e^{-jmfz} M(s)} = \frac{d_1}{(a_1 + jmf)} + \frac{d_2}{(b_1 + jf)} = G(s)$$

$$\text{note that } g(t) = d_1 e^{-a_1 t} u(t) + m d_2 e^{-mb_1 t} u(t)$$

$$P(s) = -nk_f (M(s-z) \otimes g(t))$$

Note that under the condition that  $\text{Re } s > \sigma_c$ , it follows that  $(a_1 > mb_1)$

(29)



$$\frac{F_1(f)}{2\pi k_f M C f e^{j\omega f z}} \cong \frac{-\phi_e(f)}{2\pi k_f M C f e^{j\omega f z}}$$

$$= \frac{d}{a_1 + j\omega f} + \frac{d_2}{b_1 + j\omega f} \quad \text{--- (34)}$$

Also from (2) we have

$$\frac{v(f)}{2\pi k_f M C f e^{j\omega f z}} = \frac{A_2 A_V}{2} \left(1 + \frac{c}{j\omega f}\right) \left[ \frac{d}{a_1 + j\omega f} + \frac{d_2}{b_1 + j\omega f} \right]$$

--- (35)

Let us consider a sinusoidal input  $v(t) = \begin{cases} A_m \cos \omega_f m(t - T_p), & t \geq T_p \\ 0, & t < T_p \end{cases}$

using (3) with the initial condition

$$\phi_e(t) = 0 \quad (t = T_p + \epsilon) \text{ and } \text{---}$$

investing the Fourier transform in (34) we get

$$\phi_e(t) = -2\pi k_f \left[ \frac{A_m d_1}{(a_1^2 + (\omega_f m)^2)} \left( a_1 \cos \omega_f m(t - T_p) + \omega_f m \sin(\omega_f m(t - T_p)) - a_1 e^{-a_1(t - T_p)} \right) \right.$$

$$\left. + \frac{A_m m d_2}{((\omega_f m)^2 + (\omega_f m)^2)} \left( m b_1 \cos \omega_f m(t - T_p) + \omega_f m \sin(\omega_f m(t - T_p)) - m b_1 e^{-m b_1(t - T_p)} \right) \right]$$

--- (35)

It can be shown that  $| \phi_e(t) | \ll 1$  for all  $t \geq T_p$  if  $f_m \ll \sqrt{\frac{k_v c A_2 A_V C}{2}}$  and  $k_v c A_2 A_V \gg (2 k_f m + \omega_f m)$

From eqn. (35) we have

$$\begin{aligned}
-\phi_e(t) &= \frac{2\pi k_f A_m d_1}{a_1} \left\{ \frac{a_1^2 \cos \omega_f m (t - z - T_p)}{a_1^2 + (\omega_f m)^2} \right. \\
&\quad + \frac{a_1 \omega_f m \sin \omega_f m (t - z - T_p)}{a_1^2 + (\omega_f m)^2} \\
&\quad \left. - \frac{a_1^2 e^{-\omega_f (t - z - T_p)}}{a_1^2 + (\omega_f m)^2} \right\} \\
&+ \frac{2\pi k_f (2m d_2) A_m}{(2mb_1)} \left\{ \frac{(2mb_1)^2 \cos \omega_f m (t - z - T_p)}{(2mb_1)^2 + (\omega_f m)^2} \right. \\
&\quad + \frac{(2mb_1) (\omega_f m) \sin \omega_f m (t - z - T_p)}{(2mb_1)^2 + (\omega_f m)^2} \\
&\quad \left. - \frac{(2mb_1)^2 e^{-\omega_f (t - z - T_p)}}{(2mb_1)^2 + (\omega_f m)^2} \right\}
\end{aligned}$$

since for any two positive real numbers  $x$  and  $y$   $\frac{x^2}{x^2+y^2} < 1$  and  $\frac{(xy)^2}{x^2+y^2} < \frac{1}{2}$ .

we have

$$|\phi_e(t)| < \left( 2\pi k_f A_m \left| \frac{d_1}{a_1} \right| + 2\pi k_f A_m \left| \frac{d_2}{b_1} \right| \right) \left( \frac{5}{2} \right).$$

since  $\frac{d_1}{a_1} + \frac{d_2}{b_1} = 0$ , we further get

$$\begin{aligned}
|\phi_e(t)| &< 2\pi k_f A_m \left| \frac{d_1}{a_1} \right| 5 \\
&= \frac{5 \cdot 2\pi k_f A_m}{|a_1 - 2mb_1|} = \frac{5 \cdot 2\pi k_f A_m}{(a_1 - 2mb_1)} \quad \left( \text{since } a_1 > 2mb_1 \right. \\
&\quad \left. \text{(see page 13 bottom)} \right)
\end{aligned}$$



∴ it follows that

$|p_e(s)|$  will be guaranteed

to be small if

$$(a_1 - 2nb_1) \gg mkyAm$$

$$\text{or } a_1 \gg (mkyAm + 2nb_1)$$

i.e., from the definition of  $a_1$  and  $b_1$ ,

in (27), a further sufficiency condition for  $|p_e(s)| \ll 1$  (small)

is that

$$\frac{1}{2} k_{vco} A_c A_v \gg (mkyAm + 4nC)$$

$$\text{i.e., } k_{vco} A_c A_v \gg (4kyAm + 8C)$$

This condition is required for the case when  $k_{vco} A_c A_v \gg 8C$ .

Next, we consider the more interesting case when

$$8C \gg k_{vco} A_c A_v.$$



$$\phi_e(s) = \frac{-k_f M(s) e^{-j\omega t} jf}{(jf)^2 + \left(\frac{k_{vco} A_c A_v C}{2}\right) + j \left(\frac{k_{vco} A_c A_v}{2}\right) f}$$

$$= \frac{-k_f M(s) e^{-j\omega t} jf}{\left[ \left( jf + \frac{k_{vco} A_c A_v}{4} \right)^2 + \left( \frac{k_{vco} A_c A_v C}{2} - \frac{(k_{vco} A_c A_v)^2}{16} \right) \right]}$$

we had earlier considered the case when

$$\frac{k_{vco} A_c A_v C}{2} < \frac{(k_{vco} A_c A_v)^2}{16}$$

i.e.,  $8C < k_{vco} A_c A_v$ .

Now we consider the case when

$$k_{vco} A_c A_v < 8C, \text{ in which case.}$$

we set

$$\phi_e(s) = \frac{-k_f M(s) e^{-j\omega t} jf}{\left[ \left( jf + \frac{k_{vco} A_c A_v}{4} \right)^2 - (jz_p)^2 \right]}$$

where  $z_p = \sqrt{\frac{k_{vco} A_c A_v C}{2} - \frac{(k_{vco} A_c A_v)^2}{16}} > 0$ .

$$\therefore \phi_e(s) = \frac{-k_f M(s) e^{-j\omega t} jf}{\left[ \left( j(f+z_p) + \frac{k_{vco} A_c A_v}{4} \right) \left( j(f-z_p) + \frac{k_{vco} A_c A_v}{4} \right) \right]}$$

using partial fractions we get.

$$\frac{-\phi_c(s)}{k_f m(s) e^{-j\omega_f t}} = \frac{\left(\frac{1}{2} - j \frac{k_{vco} A_c A_v}{8\eta_f}\right)}{\left(jf + \left(j\eta_f + \frac{k_{vco} A_c A_v}{4}\right)\right)} + \frac{\left(\frac{1}{2} + j \frac{k_{vco} A_c A_v}{8\eta_f}\right)}{\left(jf + \left(\frac{k_{vco} A_c A_v}{4} - j\eta_f\right)\right)}$$

corresponds to an LTI system with impulse response

$$g(t) = \left( \cos \omega_f t + \frac{n k_{vco} A_c A_v}{2\eta_f} \sin 2\omega_f t \right) e^{-\frac{n k_{vco} A_c A_v}{2} t} u(t)$$

$$\therefore \phi_c(t) = -k_f (m(t-r) \otimes g(t))$$

↑  
convolution.

Note that  $g(t)$  can also be expressed as

$$g(t) = \sqrt{4\eta_f^2 + \frac{n^2 (k_{vco} A_c A_v)^2}{4\eta_f^2}} \cos(\omega_f t - \delta) e^{-\frac{n k_{vco} A_c A_v}{2} t} u(t)$$

where  $\delta \cong \tan^{-1}\left(\frac{k_{vco} A_c A_v}{4\eta_f}\right)$ .

We note that  $g(t)$  is essentially a damped sinusoid of frequency  $\omega$ .

Since  $m(t)$  is a sinusoid of frequency  $\omega_m$ , and

$\phi_e(t) = -k_f (m(t-\tau) \otimes g(t))$  it is

clear that  $\phi_e(t)$  is small if  $\omega_m \ll \omega$

i.e.,  $\omega_m \ll \sqrt{\frac{k_{vco} A_c A_v C}{2} - \left(\frac{k_{vco} A_c A_v}{16}\right)^2}$

Exact expression for  $\phi_e(t)$  ( $t \geq \tau + T_p$ )

is given by

$$\begin{aligned} \phi_e(t) &= -k_f \int_0^{\infty} g(x) m(t-\tau-x) dx \Big|_{\substack{t-\tau-x \\ \rightarrow T_p \\ x < t-\tau}} \\ &= -k_f \int_0^{t-\tau-T_p} g(x) m(t-\tau-x) dx \\ &= -k_f A_m \int_0^{t-\tau-T_p} \cos\left(m\omega \frac{x}{2} - \frac{\pi}{4}\right) \cos m\omega (t-\tau-x - T_p) \\ &\quad e^{-\frac{\eta k_{vco} A_c A_v}{2} x} \frac{dx}{\sqrt{4\eta^2 + \frac{\eta^2}{4\omega^2} (k_{vco} A_c A_v)^2}} \end{aligned}$$



$$\frac{-\phi_e(s)}{2\pi k_f A_m} = \sqrt{1 + \left(\frac{k_{vco} A_c A_v}{4\eta_f}\right)^2} \int_0^{t-\tau-\tau_p} \cos(\omega_n \tau x - \delta) \cos(\omega_m f_m (t-\tau x - \tau_p)) e^{-\frac{\eta k_{vco} A_c A_v x}{2}} dx$$

$$= \frac{\sqrt{1 + \left(\frac{k_{vco} A_c A_v}{4\eta_f}\right)^2}}{2} \int_0^{(t-\tau-\tau_p)} \left\{ \begin{aligned} &\cos(\omega_n x (\tau_f + f_m) - \delta) \\ &\quad - \omega_n f_m t + \omega_n f_m \tau \\ &\quad + \omega_n f_m \tau_p \end{aligned} \right. + \left. \begin{aligned} &\cos(\omega_n x (\tau_f - f_m) - \delta) \\ &\quad + \omega_n f_m t - \omega_n f_m \tau \\ &\quad - \omega_n f_m \tau_p \end{aligned} \right\} e^{-\frac{\eta k_{vco} A_c A_v x}{2}} dx$$

from where it can be shown that a sufficient condition which guarantees that  $|\phi_e(s)| \ll 1$  for all  $t > \tau + \tau_p$  is

$$\frac{2\pi k_f A_m \sqrt{1 + \left(\frac{k_{vco} A_c A_v}{4\eta_f}\right)^2}}{(k_{vco} A_c A_v)} \ll 1$$

In the regime where  $C \gg \frac{k_{vco} A_c A_v}{8}$  this condition reduces to

$$k_{vco} A_c A_v \gg 2\pi k_f A_m$$

(21)

Comparison between the  
first and second order PLL's.

### PLL HOLD-IN FREQUENCY RANGE:

This refers to the maximum deviation in the frequency of the input signal from the carrier frequency, such that the PLL can still lock (i.e.,  $\phi_e(t)$  is small).

### FIRST ORDER-PLL:

We required that

$$\Delta f_{\max} = k_f A_m \ll \frac{k_{vco} A_c A_v}{2}$$

### SECOND ORDER PLL:

For a second order PLL with

$$\zeta \gg \frac{k_{vco} A_c A_v}{8}, \text{ we required that}$$

$$\Delta f_{\max} = k_f A_m \ll \frac{k_{vco} A_c A_v}{2}$$

i.e., in both the first & second order PLL's  $k_{vco}$  can be used to control the PLL HOLD-IN FREQUENCY RANGE.

LOOP BANDWIDTH

We define loop bandwidth to be the maximum possible bandwidth of  $m(t)$  such that  $v(t)$  is a distortion-free scalar multiple of  $m(t-z)$ .

Consider an input  $m(t)$  consisting of two sinusoids

$$m(t) = \begin{cases} A_m \cos 2\pi f_{m1}(t - T_p) + A_m \cos 2\pi f_{m2}(t - T_p) & , t \geq T_p \\ 0 & , t < T_p \end{cases}$$

then using (16), the output  $v(t)$  is given by

$$v(t) \approx A_m \left\{ G_1 \cos \left( 2\pi f_{m1}(t - z - T_p) - \tan^{-1} \frac{2\pi f_{m1}}{a} \right) + G_2 \cos \left( 2\pi f_{m2}(t - z - T_p) - \tan^{-1} \left( \frac{2\pi f_{m2}}{a} \right) \right) \right\}$$

Consider an example where

$$f_{m1} = \frac{a}{2\pi} \text{ and } f_{m2} = \frac{a}{2\pi} (\sqrt{3})$$

then since  $G(f) = \frac{a}{\sqrt{a^2 + (2\pi f)^2}}$ ,  $G(f_{m1}) = \frac{1}{\sqrt{2}}$  and  $G(f_{m2}) = \frac{1}{2}$



∴ the output

$$v(t) \approx \frac{k_f A_m}{2 k_v \omega} \left( \sqrt{2} \cos(m f \omega_1 (t - 2T_p) - \pi/4) + \cos(m f \omega_2 (t - 2T_p) - \pi/3) \right)$$

which is a distorted version of  $m(t - 2)$  due to the different phase shifts i.e.,  $(-\pi/4$  and  $-\pi/3)$  in the constituent sinusoids of  $m(t)$  (Draw this in Matlab and see the distortion produced)

For the first order PLL loop

$$G(f) = \frac{a}{\sqrt{a^2 + \omega^2}}, \text{ and } \angle(f) = \tan^{-1}\left(\frac{\omega}{a}\right)$$

∴ the only way to control distortion in the first order loop is to ensure that  $G(f) \approx 1$  and  $\angle(f) \propto f$  (i.e., linear phase)

Given  $G(f)$  and  $\angle(f)$  (above), it is clear that  $G(f) \approx 1$  and  $\angle(f) \propto f$  if and only if  $|f| \ll \left(\frac{a}{\omega}\right)$ .

therefore for the first order loop (24)

the loop bandwidth (BL) must satisfy

$$BL \ll \frac{\omega}{n} = \frac{k_{vco} A_c A_v}{2}$$

∴ The  $n$  in the first order PLL the same parameter  $k_{vco} A_c A_v$  controls both the loop bandwidth as well as the PLL lock-in frequency range.

For a second order PLL, we have (using equation (25))

$$\frac{V(f)}{M(f)e^{j2\pi f t}} = \frac{n k_f A_c A_v (c + jf)}{(n k_{vco} A_c A_v c + j n k_{vco} A_c A_v f - 2\pi f^2)}$$

from where it is clear that for  $f$  such that

$$n k_{vco} A_c A_v c \gg 2\pi f^2$$

i.e.  $f \ll \sqrt{\frac{k_{vco} A_c A_v c}{2}}$  we have -

$$\begin{aligned} \frac{V(f)}{M(f)e^{j2\pi f t}} &\approx \frac{n k_f A_c A_v (c + jf)}{(n k_{vco} A_c A_v c + j n k_{vco} A_c A_v f)} \\ &= \frac{k_f}{k_{vco}} \left( f \ll \sqrt{\frac{k_{vco} A_c A_v c}{2}} \right) \end{aligned}$$

In a second order loop if  $m(t)$  has its highest frequency component  $\ll \sqrt{\frac{k_{vco} A_c A_v C}{2}}$

then  $v(t) \approx \frac{k_f}{k_{vco}} m(t-\tau)$ , i.e., the PLL output  $\uparrow v(t)$  will be a distortion-free delayed version of  $m(t)$ .

DISCUSSION (SUMMARY)

In a first order PLL, both the PLL lock-in frequency range and the loop bandwidth are controlled by the same parameter  $\frac{k_{vco} A_c A_v}{2}$ . However the gain of the PLL is  $\frac{k_f}{k_{vco}}$ , i.e.,

a large  $k_{vco}$  will imply a loss in the amplitude at the PLL output. There is therefore a trade-off between a higher loop bandwidth/PLL lock-in frequency range and the PLL gain (in the case of a first order loop). However in the case of a second-order PLL, by choosing  $C \gg \frac{k_{vco} A_c A_v}{8}$ , we can control the PLL lock-in frequency range using  $\frac{k_{vco} A_c A_v}{2}$ , and independently control the loop bandwidth using  $C$  (since loop bandwidth  $= 0.1 \sqrt{\frac{k_{vco} A_c A_v C}{2}}$ )



∴ if we want to have a larger PLL loop bandwidth without increasing  $\frac{K_{VCO} A_2 A_V}{2}$ , then we can do so in a second-order PLL using the control parameter  $C$ .

	First-order PLL	Second-order PLL
PLL HOLD-IN FREQUENCY	$\ll \frac{K_{VCO} A_2 A_V}{2}$	$\ll \frac{K_{VCO} A_2 A_V}{2}$
PLL LOOP BANDWIDTH	$\ll \frac{K_{VCO} A_2 A_V}{2}$	$\ll \sqrt{\frac{K_{VCO} A_2 A_V C}{2}}$ (assumes $C \gg \frac{K_{VCO} A_2 A_V}{8}$ )