

## Entropy and Temperature

$$dU = \delta q_{rev} + \delta w_{rev} = TdS - PdV$$

We can get,

$$\left(\frac{\partial S}{\partial T}\right)_V = \frac{C_V}{T} \text{ and } \left(\frac{\partial S}{\partial V}\right)_T = \frac{1}{T} \left[ P + \left(\frac{\partial U}{\partial V}\right)_T \right]$$

$$\Delta S = S(T_2) - S(T_1) = \int_{T_1}^{T_2} \frac{C_V(T)dT}{T}$$

More often than not we deal with constant pressure processes.

$$dH = d(U + PV) = dU + PdV + VdP = TdS + VdP$$

$$\left(\frac{\partial S}{\partial T}\right)_P = \frac{C_P}{T} \text{ and } \left(\frac{\partial S}{\partial P}\right)_T = \frac{1}{T} \left[ \left(\frac{\partial H}{\partial P}\right)_T - V \right]$$

$$\Delta S = S(T_2) - S(T_1) = \int_{T_1}^{T_2} \frac{C_P(T)dT}{T}$$

If we let  $T_1 = 0$

$$S(T) = S(0 \text{ K}) + \int_0^T \frac{C_P(T')dT'}{T'}$$

We can therefore calculate the entropy of a substance if we know  $S(0)$  and  $C_P(T)$

## Third Law of Thermodynamics

A perfectly crystalline substance has zero entropy at absolute zero

$$S = k_B \ln W = -k_B \sum_j p_j \ln p_j$$

Since at 0 K there is no thermal energy, only the G.S. is occupied, i.e.  $p_0 = 1$  and all other states are vacant, or  $p_j = 0$  when  $j > 0$ .  $\therefore S = 0$ .

In case there is degeneracy of the G.S. i.e.  $n$  available states of energy  $E_0$  then each of them has a probability  $\frac{1}{n}$  to be filled.

$$S(0\text{ K}) = -k_B \sum_{j=1}^n \frac{1}{n} \ln \frac{1}{n} = k_B \ln n \quad \dots \text{a very small number even for large value of } n$$

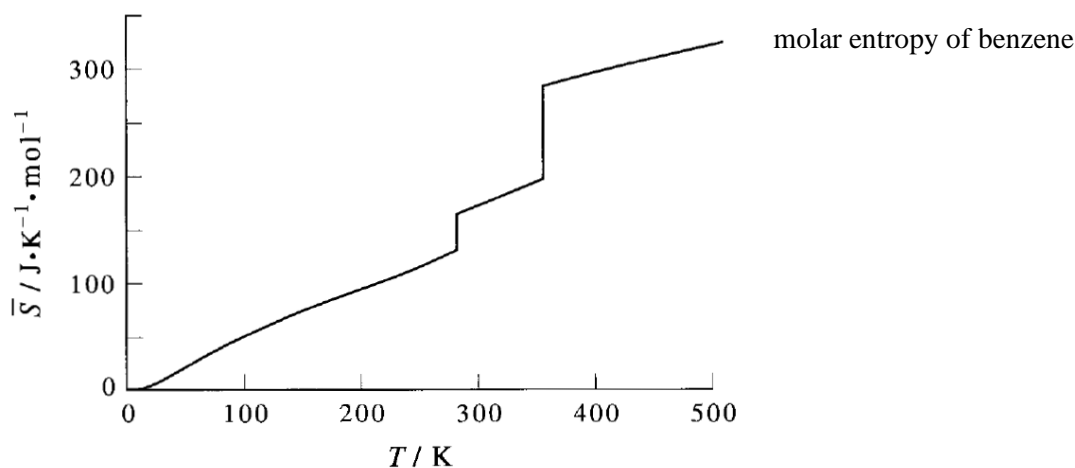
Now we can write,

$$S(T) = \int_0^T \frac{C_P(T') dT'}{T'}$$

For calculating we need to take care of phase transitions

$$\Delta_{trs} S = \frac{q_{rev}}{T} = \frac{\Delta_{trs} H}{T_{trs}}$$

$$S(T) = \int_0^{T_{fus}} \frac{C_P^s(T') dT'}{T'} + \frac{\Delta_{fus} H}{T_{fus}} + \int_{T_{fus}}^{T_{vap}} \frac{C_P^l(T') dT'}{T'} + \frac{\Delta_{vap} H}{T_{vap}} + \int_{T_{vap}}^T \frac{C_P^g(T') dT'}{T'}$$



## Contributions to the Absolute Entropy

$$S = k_B \ln Q + k_B T \left( \frac{\partial \ln Q}{\partial T} \right)_{N,V} = k_B \ln \frac{q^N}{N!} + k_B T \left( \frac{\partial \ln(q^N/N!)}{\partial T} \right)_V$$

$$= Nk_B \ln q - k_B \ln N! + Nk_B T \left( \frac{\partial \ln q}{\partial T} \right)_V$$

Using Stirling's Approximation ( $\ln N! = N \ln N - N$ ):  $S = Nk_B + Nk_B \ln \frac{q}{N} + Nk_B T \left( \frac{\partial \ln q}{\partial T} \right)_V$

As an example, take the diatomic ideal gas.  $N_2(g)$  at 298.15 K

$$\frac{\bar{S}}{R} = \ln \left[ \left( \frac{2\pi M k_B T}{h^2} \right)^{3/2} \frac{\bar{V} e^{5/2}}{N_A} \right] + \ln \frac{T}{2\Theta_{rot}} - \ln(1 - e^{-\Theta_{vib}/T}) + \frac{\Theta_{vib}/T}{e^{-\Theta_{vib}/T} - 1} + \ln g_{e1}$$

$$\text{Using, } \Theta_{rot} = 2.88 \text{ K}, \Theta_{vib} = 3374 \text{ K}, g_{e1} = 1$$

$$\text{Standard entropy, } S^0(298.15 \text{ K}) = S_{trans}^0 + S_{rot}^0 + S_{vib}^0 + S_{elec}^0$$

$$= 150.4 + 41.13 + 1.15 \times 10^{-3} + 0 = 191.5 \text{ JK}^{-1}\text{mol}^{-1}$$

From heat capacity:  $191.6 \text{ JK}^{-1}\text{mol}^{-1}$

$$S(T') = \int_0^{T_{fus}} \frac{C_p^s(T) dT}{T} + \frac{\Delta_{fus}H}{T_{fus}} + \int_{T_{fus}}^{T_{vap}} \frac{C_p^l(T) dT}{T} + \frac{\Delta_{vap}H}{T_{vap}} + \int_{T_{vap}}^{T'} \frac{C_p^g(T) dT}{T}$$

## Molecular Partition Functions

Monatomic ideal gas

$$q(V, T) = \left( \frac{2\pi m k_B T}{h^2} \right)^{3/2} V \cdot g_{e1}$$

Diatomic ideal gas

$$q(V, T) = \left( \frac{2\pi M k_B T}{h^2} \right)^{3/2} V \cdot \frac{T}{\sigma \Theta_{rot}} \cdot \frac{e^{-\Theta_{vib}/2T}}{1 - e^{-\Theta_{vib}/T}} \cdot g_{e1} e^{D_e/k_B T}$$

Linear polyatomic ideal gas

$$q(V, T) = \left( \frac{2\pi M k_B T}{h^2} \right)^{3/2} V \cdot \frac{T}{\sigma \Theta_{rot}} \cdot \left[ \prod_{j=1}^{3n-5} \frac{e^{-\Theta_{vib,j}/2T}}{1 - e^{-\Theta_{vib,j}/T}} \right] \cdot g_{e1} e^{D_e/k_B T}$$

Non-linear polyatomic ideal gas

$$q(V, T) = \left( \frac{2\pi M k_B T}{h^2} \right)^{3/2} V \cdot \frac{\pi^{1/2}}{\sigma} \left( \frac{T^3}{\Theta_A \Theta_B \Theta_C} \right)^{1/2} \cdot \left[ \prod_{j=1}^{3n-6} \frac{e^{-\Theta_{vib,j}/2T}}{1 - e^{-\Theta_{vib,j}/T}} \right] \cdot g_{e1} e^{D_e/k_B T}$$

### Some useful relations

$$dU = \delta q_{rev} + \delta w_{rev}$$

$$dU = TdS - PdV$$

$$\left(\frac{\partial U}{\partial T}\right)_V = T \left(\frac{\partial S}{\partial T}\right)_V$$

$$\left(\frac{\partial S}{\partial T}\right)_V = \frac{C_V(T)}{T}$$

$$dH = d(U + PV) = dU + PdV + VdP = TdS + VdP$$

$$\left(\frac{\partial S}{\partial T}\right)_P = \frac{C_P(T)}{T}$$